

ECE276A: Sensing & Estimation in Robotics

Lecture 12: $SO(3)$ and $SE(3)$ Geometry and Kinematics

Instructor:

Nikolay Atanasov: natanasov@ucsd.edu

Teaching Assistants:

Qiaojun Feng: qif007@eng.ucsd.edu

Arash Asgharivaskasi: aasghari@eng.ucsd.edu

Thai Duong: tduong@eng.ucsd.edu

Yiran Xu: y5xu@eng.ucsd.edu

UC San Diego

JACOBS SCHOOL OF ENGINEERING
Electrical and Computer Engineering

Representations of Orientation

- ▶ **Rotation Matrix:** an element of the **Special Orthogonal Group**:

$$R \in SO(3) := \left\{ R \in \mathbb{R}^{3 \times 3} \mid \underbrace{R^T R = I}_{\text{distances preserved}}, \underbrace{\det(R) = 1}_{\text{no reflection}} \right\}$$

- ▶ **Euler Angles:** roll ϕ , pitch θ , roll ψ specifying a **rzyx** rotation:

$$R = R_z(\psi)R_y(\theta)R_x(\phi)$$

- ▶ **Axis-Angle:** $\boldsymbol{\theta} \in \mathbb{R}^3$ specifying a rotation about an axis $\boldsymbol{\eta} := \frac{\boldsymbol{\theta}}{\|\boldsymbol{\theta}\|}$ through an angle $\theta := \|\boldsymbol{\theta}\|$:

$$R = \exp(\hat{\boldsymbol{\theta}}) = I + \hat{\boldsymbol{\theta}} + \frac{1}{2!}\hat{\boldsymbol{\theta}}^2 + \frac{1}{3!}\hat{\boldsymbol{\theta}}^3 + \dots$$

- ▶ **Unit Quaternion:** $\mathbf{q} = [q_s, \mathbf{q}_v] \in \{q \in \mathbb{H} \mid q_s^2 + \mathbf{q}_v^T \mathbf{q}_v = 1\}$:

$$R = E(\mathbf{q})G(\mathbf{q})^T \quad \begin{aligned} E(\mathbf{q}) &= [-\mathbf{q}_v, q_s I + \hat{\mathbf{q}}_v] \\ G(\mathbf{q}) &= [-\mathbf{q}_v, q_s I - \hat{\mathbf{q}}_v] \end{aligned}$$

Special Orthogonal Group $SO(3)$

- ▶ The orientation R of a rigid body can be described by a matrix in the **special orthogonal group**:

$$SO(3) := \left\{ R \in \mathbb{R}^{3 \times 3} \mid R^T R = I, \det(R) = 1 \right\}$$

- ▶ It can be verified that $SO(3)$ satisfies all requirements of a group:
 - ▶ **Closure**: $R_1 R_2 \in SO(3)$
 - ▶ **Identity**: $I \in SO(3)$
 - ▶ **Inverse**: $R^{-1} = R^T \in SO(3)$
 - ▶ **Associativity**: $(R_1 R_2) R_3 = R_1 (R_2 R_3)$ for all $R_1, R_2, R_3 \in SO(3)$

Special Euclidean Group $SE(3)$

- ▶ The pose T of a rigid body can be described by a matrix in the **special Euclidean group**:

$$SE(3) := \left\{ T := \begin{bmatrix} R & \mathbf{p} \\ \mathbf{0}^\top & 1 \end{bmatrix} \mid R \in SO(3), \mathbf{p} \in \mathbb{R}^3 \right\} \subset \mathbb{R}^{4 \times 4}$$

- ▶ It can be verified that $SE(3)$ satisfies all requirements of a group:
 - ▶ **Closure:** $T_1 T_2 = \begin{bmatrix} R_1 & \mathbf{p}_1 \\ \mathbf{0}^\top & 1 \end{bmatrix} \begin{bmatrix} R_2 & \mathbf{p}_2 \\ \mathbf{0}^\top & 1 \end{bmatrix} = \begin{bmatrix} R_1 R_2 & R_1 \mathbf{p}_2 + \mathbf{p}_1 \\ \mathbf{0}^\top & 1 \end{bmatrix} \in SE(3)$
 - ▶ **Identity:** $\begin{bmatrix} I & \mathbf{0} \\ \mathbf{0}^\top & 1 \end{bmatrix} \in SE(3)$
 - ▶ **Inverse:** $\begin{bmatrix} R & \mathbf{p} \\ \mathbf{0}^\top & 1 \end{bmatrix}^{-1} = \begin{bmatrix} R^\top & -R^\top \mathbf{p} \\ \mathbf{0}^\top & 1 \end{bmatrix} \in SE(3)$
 - ▶ **Associativity:** $(T_1 T_2) T_3 = T_1 (T_2 T_3)$ for all $T_1, T_2, T_3 \in SE(3)$

Matrix Lie Group

- ▶ $SO(3)$ and $SE(3)$ are **matrix Lie groups**
- ▶ A **group** is a set of elements with an operation that combines any two elements to form a third one also in the set. A group satisfies four axioms: closure, associativity, identity, and invertibility
- ▶ A **Lie group** is a group that is also a differentiable manifold with the property that the group operations are smooth
- ▶ A **matrix** Lie group further specifies that the group elements are matrices, the combination operation is matrix multiplication, and the inversion operation is matrix inversion
- ▶ The **exponential map** relates a matrix Lie group to its **Lie algebra**

$$\exp(A) = \sum_{n=0}^{\infty} \frac{1}{n!} A^n \qquad \log(A) = \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{n} (A - I)^n$$

Lie Algebra

- ▶ A **Lie algebra** is associated with every matrix Lie group.
- ▶ A Lie algebra is a vector space \mathbb{V} over some field \mathbb{F} with a binary operation, $[\cdot, \cdot]$, called a **Lie bracket**
- ▶ The vector space of a Lie algebra is the **tangent space** of the associated Lie group at the identity element of the group
- ▶ For all $X, Y, Z \in \mathbb{V}$ and $a, b \in \mathbb{F}$, the Lie bracket satisfies:

$$\text{closure : } [X, Y] \in \mathbb{V}$$

$$\text{bilinearity : } [aX + bY, Z] = a[X, Z] + b[Y, Z]$$

$$[Z, aX + bY] = a[Z, X] + b[Z, Y]$$

$$\text{alternating : } [X, X] = 0$$

$$\text{Jacobi identity : } [X, [Y, Z]] + [Y, [Z, X]] + [Z, [X, Y]] = 0$$

Lie Group and Lie Algebra Visualization

- ▶ **Lie Group:** free of singularities but has constraints
- ▶ **Lie Algebra:** free of constraints but has singularities

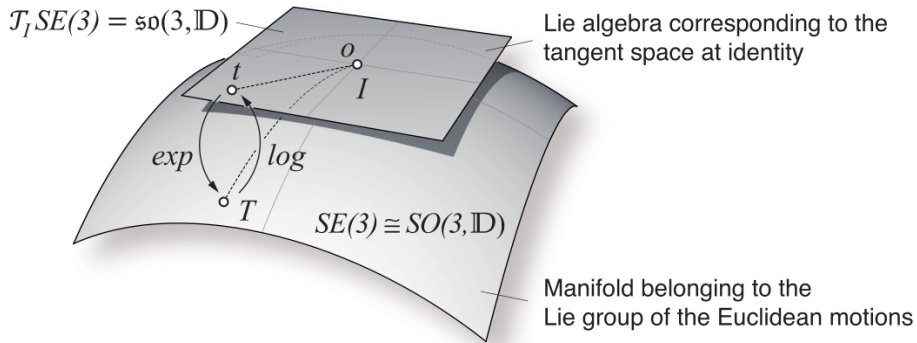


Figure: $SE(3)$ and the corresponding Lie algebra $\mathfrak{se}(3)$ as tangent space at identity

- ▶ The elements of matrix Lie groups do not satisfy some basic operations that we normally take for granted

$SO(3)$ Geometry

Special Orthogonal Lie Algebra $\mathfrak{so}(3)$

- ▶ The Lie algebra of $SO(3)$ is the space of skew-symmetric matrices

$$\mathfrak{so}(3) := \{\hat{\theta} \in \mathbb{R}^{3 \times 3} \mid \theta \in \mathbb{R}^3\}$$

- ▶ The **Lie bracket** of $\mathfrak{so}(3)$ is:

$$[\hat{\theta}_1, \hat{\theta}_2] = \hat{\theta}_1 \hat{\theta}_2 - \hat{\theta}_2 \hat{\theta}_1 = (\hat{\theta}_1 \theta_2)^\wedge \in \mathfrak{so}(3)$$

- ▶ **Generators of $\mathfrak{so}(3)$:** derivatives of rotations around each standard axis:

$$G_x = \left. \frac{d}{d\phi} R_x(\phi) \right|_{\phi=0} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{bmatrix} \quad G_y = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{bmatrix} \quad G_z = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

- ▶ The elements $\hat{\theta} = \alpha_1 G_x + \alpha_2 G_y + \alpha_3 G_z \in \mathfrak{so}(3)$ are linear combinations of generators and can be mapped to $SO(3)$ via the **exponential map**:

$$R = \exp(\hat{\theta}) = I + \hat{\theta} + \frac{1}{2!} \hat{\theta}^2 + \frac{1}{3!} \hat{\theta}^3 + \dots \quad \theta = \log(R)^\vee$$

Exponential Map from $\mathfrak{so}(3)$ to $SO(3)$

- ▶ The exponential map is **surjective** but **not injective**, i.e., every element of $SO(3)$ can be generated from multiple elements of $\mathfrak{so}(3)$
- ▶ **Rodrigues Formula**: a closed-form expression for the exponential map from $\mathfrak{so}(3)$ to $SO(3)$ using that $\hat{\theta}^{2n+1} = (-\theta^\top \theta)^n \hat{\theta}$:

$$\begin{aligned} R = \exp(\hat{\theta}) &= I + \sum_{n=1}^{\infty} \frac{1}{n!} \hat{\theta}^n = I + \sum_{n=0}^{\infty} \frac{1}{(2n+1)!} \hat{\theta}^{2n+1} + \sum_{n=0}^{\infty} \frac{1}{(2n+2)!} \hat{\theta}^{2n+2} \\ &= I + \left(\sum_{n=0}^{\infty} \frac{(-1)^n \|\theta\|^{2n}}{(2n+1)!} \right) \hat{\theta} + \left(\sum_{n=0}^{\infty} \frac{(-1)^n \|\theta\|^{2n}}{(2n+2)!} \right) \hat{\theta}^2 \\ &= I + \left(\frac{\sin \|\theta\|}{\|\theta\|} \right) \hat{\theta} + \left(\frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta}^2 \end{aligned}$$

- ▶ Any vector $\theta + 2\pi k$ for integer k leads to the same $R \in SO(3)$
- ▶ The exponential map is **not commutative**, $e^{\hat{\theta}_1} e^{\hat{\theta}_2} \neq e^{\hat{\theta}_2} e^{\hat{\theta}_1} \neq e^{\hat{\theta}_1 + \hat{\theta}_2}$, unless $\hat{\theta}_1 \hat{\theta}_2 = \hat{\theta}_2 \hat{\theta}_1$, i.e., the **Lie bracket** on $\mathfrak{so}(3)$, $[\hat{\theta}_1, \hat{\theta}_2] = 0$.

Logarithm Map from $SO(3)$ to $\mathfrak{so}(3)$

- ▶ $\forall R \in SO(3)$, there exists a (not unique) $\theta \in \mathbb{R}^3$ such that $R = \exp(\hat{\theta})$
- ▶ The **logarithm map** $\log : SO(3) \rightarrow \mathfrak{so}(3)$ is the inverse of $\exp(\hat{\theta})$:

$$\theta = \|\theta\| = \arccos\left(\frac{\text{tr}(R) - 1}{2}\right)$$

$$\eta = \frac{\theta}{\|\theta\|} = \frac{1}{2 \sin(\|\theta\|)} \begin{bmatrix} R_{32} - R_{23} \\ R_{13} - R_{31} \\ R_{21} - R_{12} \end{bmatrix}$$

$$\hat{\theta} = \log(R) = \frac{\|\theta\|}{2 \sin \|\theta\|} (R - R^T)$$

- ▶ If $R = I$, then $\theta = 0$ and η is undefined

- ▶ If $\text{tr}(R) = -1$, then $\theta = \pi$ and for any $i \in \{1, 2, 3\}$:

$$\eta = \frac{1}{\sqrt{2(1 + e_i^T R e_i)}} (I + R) e_i$$

- ▶ The log map has a singularity at $\theta = 0$ because there are infinite choices of rotation axes or equivalently the exponential map is many-to-one.
- ▶ The matrix exponential “integrates” $\hat{\theta} \in \mathfrak{se}(3)$ for one second; the matrix logarithm “differentiates” $R \in SO(3)$ to obtain $\hat{\theta} \in \mathfrak{se}(3)$

$SO(3)$ Jacobians

- ▶ The **left Jacobian** of $SO(3)$ is the matrix:

$$J_L(\boldsymbol{\theta}) := \sum_{n=0}^{\infty} \frac{1}{(n+1)!} (\hat{\boldsymbol{\theta}})^n \quad R = I + \hat{\boldsymbol{\theta}} J_L(\boldsymbol{\theta})$$

- ▶ The **right Jacobian** of $SO(3)$ is the matrix:

$$J_R(\boldsymbol{\theta}) := \sum_{n=0}^{\infty} \frac{1}{(n+1)!} (-\hat{\boldsymbol{\theta}})^n \quad J_R(\boldsymbol{\theta}) = J_L(-\boldsymbol{\theta}) = J_L(\boldsymbol{\theta})^T = R^T J_L(\boldsymbol{\theta})$$

- ▶ **Baker-Campbell-Hausdorff Formulas:** the $SO(3)$ Jacobians relate small perturbations in $\mathfrak{so}(3)$ to small perturbations in $SO(3)$:

$$\begin{aligned} \exp((\boldsymbol{\theta} + \delta\boldsymbol{\theta})^\wedge) &\approx \exp(\hat{\boldsymbol{\theta}}) \exp((J_R(\boldsymbol{\theta})\delta\boldsymbol{\theta})^\wedge) \\ &\approx \exp((J_L(\boldsymbol{\theta})\delta\boldsymbol{\theta})^\wedge) \exp(\hat{\boldsymbol{\theta}}) \end{aligned}$$

$$\log(\exp(\hat{\boldsymbol{\theta}}_1) \exp(\hat{\boldsymbol{\theta}}_2))^\vee \approx \begin{cases} J_L(\boldsymbol{\theta}_2)^{-1}\boldsymbol{\theta}_1 + \boldsymbol{\theta}_2 & \text{if } \boldsymbol{\theta}_1 \text{ is small} \\ \boldsymbol{\theta}_1 + J_R(\boldsymbol{\theta}_1)^{-1}\boldsymbol{\theta}_2 & \text{if } \boldsymbol{\theta}_2 \text{ is small} \end{cases}$$

Closed-forms of the $SO(3)$ Jacobians

$$J_L(\theta) = I + \left(\frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta} + \left(\frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) \hat{\theta}^2 \approx I + \frac{1}{2} \hat{\theta}$$

$$J_L(\theta)^{-1} = I - \frac{1}{2} \hat{\theta} + \left(\frac{1}{\|\theta\|^2} - \frac{1 + \cos \|\theta\|}{2\|\theta\| \sin \|\theta\|} \right) \hat{\theta}^2 \approx I - \frac{1}{2} \hat{\theta}$$

$$J_R(\theta) = I - \left(\frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta} + \left(\frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) \hat{\theta}^2 \approx I - \frac{1}{2} \hat{\theta}$$

$$J_R(\theta)^{-1} = I + \frac{1}{2} \hat{\theta} + \left(\frac{1}{\|\theta\|^2} - \frac{1 + \cos \|\theta\|}{2\|\theta\| \sin \|\theta\|} \right) \hat{\theta}^2 \approx I + \frac{1}{2} \hat{\theta}$$

$$J_L(\theta) J_L(\theta)^T = I + \left(1 - 2 \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta}^2 \succ 0$$

$$\left(J_L(\theta) J_L(\theta)^T \right)^{-1} = I + \left(1 - 2 \frac{\|\theta\|^2}{1 - \cos \|\theta\|} \right) \hat{\theta}^2$$

Distances in $SO(3)$

- ▶ There are two ways to define the difference between two rotations:

$$\boldsymbol{\theta}_{12} = \log \left(R_1^\top R_2 \right)^\vee \quad \boldsymbol{\theta}_{21} = \log \left(R_2 R_1^\top \right)^\vee \quad R_1, R_2 \in SO(3)$$

- ▶ Inner product on $\mathfrak{so}(3)$:

$$\langle \hat{\boldsymbol{\theta}}_1, \hat{\boldsymbol{\theta}}_2 \rangle = \frac{1}{2} \operatorname{tr} \left(\hat{\boldsymbol{\theta}}_1 \hat{\boldsymbol{\theta}}_2^\top \right) = \boldsymbol{\theta}_1^\top \boldsymbol{\theta}_2$$

- ▶ The metric distance between two rotations may be defined in two ways as the magnitude of the rotation difference:

$$\theta_{12} := \sqrt{\langle \log \left(R_1^\top R_2 \right), \log \left(R_1^\top R_2 \right) \rangle} = \|\boldsymbol{\theta}_{12}\|$$

$$\theta_{21} := \sqrt{\langle \log \left(R_2 R_1^\top \right), \log \left(R_2 R_1^\top \right) \rangle} = \|\boldsymbol{\theta}_{21}\|$$

Integration in $SO(3)$

- ▶ The distance between a rotation $R = \exp(\hat{\theta})$ and a small perturbation $\exp((\theta + \delta\theta)^\wedge)$ can be approximated using the BCH formulas:

$$\log \left(\exp(\hat{\theta})^T \exp((\theta + \delta\theta)^\wedge) \right)^\vee \approx \log \left(R^T R \exp((J_R(\theta)\delta\theta)^\wedge) \right)^\vee = J_R(\theta)\delta\theta$$

$$\log \left(\exp((\theta + \delta\theta)^\wedge) \exp(\hat{\theta})^T \right)^\vee \approx \log \left(\exp((J_L(\theta)\delta\theta)^\wedge) R^T R \right)^\vee = J_L(\theta)\delta\theta$$

- ▶ Regardless of which distance metric we use, the infinitesimal volume element is the same:

$$\det(J_L(\theta)) = \det(J_R(\theta)) \quad dR = |\det(J(\theta))| d\theta = 2 \left(\frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) d\theta$$

- ▶ Integrating functions of rotations can then be carried out as follows:

$$\int_{SO(3)} f(R) dR = \int_{\|\theta\| < \pi} f(\theta) |\det(J(\theta))| d\theta$$

Derivatives in $SO(3)$

- ▶ Consider $\mathbf{s} \in \mathbb{R}^3$ rotated by a rotation matrix $R \in SO(3)$ to a new frame
- ▶ How do we compute the derivative of $R\mathbf{s}$ with respect to the rotation R ?
- ▶ Let $\boldsymbol{\theta} \in \mathbb{R}^3$ be the Lie algebra vector representing R , i.e., $R = \exp(\hat{\boldsymbol{\theta}})$
- ▶ We can compute derivatives with respect to the elements of $\boldsymbol{\theta}$:

$$\begin{aligned} \frac{\partial R\mathbf{s}}{\partial \theta_i} &= \lim_{h \rightarrow 0} \frac{\exp((\boldsymbol{\theta} + h\mathbf{e}_i)^\wedge) \mathbf{s} - \exp(\hat{\boldsymbol{\theta}}) \mathbf{s}}{h} \\ &\stackrel{\text{BCH}}{\text{Formula}} \lim_{h \rightarrow 0} \frac{\exp((hJ_L(\boldsymbol{\theta})\mathbf{e}_i)^\wedge) \exp(\hat{\boldsymbol{\theta}}) \mathbf{s} - \exp(\hat{\boldsymbol{\theta}}) \mathbf{s}}{h} \\ &\stackrel{\exp(\delta\hat{\boldsymbol{\theta}}) \approx I + \delta\hat{\boldsymbol{\theta}}}{=} \lim_{h \rightarrow 0} \frac{(I + h(J_L(\boldsymbol{\theta})\mathbf{e}_i)^\wedge) \exp(\hat{\boldsymbol{\theta}}) \mathbf{s} - \exp(\hat{\boldsymbol{\theta}}) \mathbf{s}}{h} \\ &= (J_L(\boldsymbol{\theta})\mathbf{e}_i)^\wedge R\mathbf{s} = -(R\mathbf{s})^\wedge J_L(\boldsymbol{\theta})\mathbf{e}_i \end{aligned}$$

- ▶ Stacking the three directional derivatives: $\frac{\partial R\mathbf{s}}{\partial \boldsymbol{\theta}} = -(R\mathbf{s})^\wedge J_L(\boldsymbol{\theta})$

Derivatives in $SO(3)$

- ▶ **Perturbation in $\mathfrak{so}(3)$:** the gradient can also be obtained via a small perturbation $\delta\theta$ to the axis-angle vector θ :

$$\begin{aligned}\exp((\theta + \delta\theta)^\wedge) \mathbf{s} &\stackrel{\text{BCH}}{\approx} \exp((J_L\theta)\delta\theta)^\wedge \exp(\hat{\theta})\mathbf{s} \\ &\approx (I + (J_L(\theta)\delta\theta)^\wedge) \exp(\hat{\theta})\mathbf{s} \\ &= R\mathbf{s} + (J_L(\theta)\delta\theta)^\wedge R\mathbf{s} = R\mathbf{s} - \underbrace{(R\mathbf{s})^\wedge J_L(\theta)}_{\frac{\partial R\mathbf{s}}{\partial \theta}} \delta\theta\end{aligned}$$

- ▶ This is the same as using first-order Taylor series to identify the Jacobian of a function $f(\mathbf{x})$:

$$f(\mathbf{x} + \delta\mathbf{x}) \approx f(\mathbf{x}) + \left[\frac{\partial f}{\partial \mathbf{x}}(\mathbf{x}) \right] \delta\mathbf{x}$$

- ▶ **Perturbation in $SO(3)$:** a small perturbation $\psi := J_L(\theta)\delta\theta$ may also be applied directly to R :

$$\exp(\hat{\psi})R\mathbf{s} \approx (I + \hat{\psi})R\mathbf{s} = R\mathbf{s} - (R\mathbf{s})^\wedge \psi$$

Gradient Descent in $SO(3)$

- ▶ Consider $\min_{\mathbf{x}} f(\mathbf{x})$
- ▶ **Gradient descent:** given an initial guess $\mathbf{x}^{(k)}$ take a step of size $\alpha^{(k)} > 0$ along a descent direction $\delta\mathbf{x}^{(k)}$ such that:

$$\mathbf{x}^{(k+1)} = \mathbf{x}^{(k)} + \alpha^{(k)}\delta\mathbf{x}^{(k)} \quad \text{and} \quad \delta\mathbf{x}^{(k)} = -\nabla f(\mathbf{x}^{(k)})$$

- ▶ Consider $\min_R f(R\mathbf{s})$
- ▶ **Gradient descent in $SO(3)$:** given an initial guess $R^{(k)}$ take a step of size $\alpha^{(k)} > 0$ along a descent direction $\psi^{(k)}$ such that:

$$R^{(k+1)} = \exp\left(\alpha^{(k)}\hat{\psi}^{(k)}\right) R^{(k)} \quad \text{and} \quad \psi^{(k)} = -\delta^{(k)}$$

where $\delta^{(k)}$ should be the gradient of f wrt R evaluated at $R^{(k)}\mathbf{s}$

Choosing a Descent Direction in $SO(3)$

- ▶ Use a perturbation $\psi^{(k)}$ around the initial guess $R^{(k)}$ to determine the gradient $\delta^{(k)}$:

$$\begin{aligned} f\left(\exp(\hat{\psi}^{(k)})R^{(k)}\mathbf{s}\right) &\approx f\left((I + \hat{\psi}^{(k)})R^{(k)}\mathbf{s}\right) \\ &= f\left(R^{(k)}\mathbf{s} - \left(R^{(k)}\mathbf{s}\right)^\wedge \psi^{(k)}\right) \\ &\approx f\left(R^{(k)}\mathbf{s}\right) - \underbrace{\nabla f\left(R^{(k)}\mathbf{s}\right)^\top \left(R^{(k)}\mathbf{s}\right)^\wedge}_{\delta^{(k)\top}} \psi^{(k)} \end{aligned}$$

- ▶ **Gradient descent in $SO(3)$:** given an initial guess $R^{(k)}$ take a step of size $\alpha^{(k)} > 0$:

$$\begin{aligned} \delta^{(k)} &= \left(R^{(k)}\mathbf{s}\right)^\wedge \nabla f\left(R^{(k)}\mathbf{s}\right) \\ R^{(k+1)} &= \exp\left(-\alpha^{(k)}\hat{\delta}^{(k)}\right) R^{(k)} \end{aligned}$$

Gauss-Newton Optimization in $SO(3)$

- ▶ Optimization problem:

$$\min_R f(R) := \frac{1}{2} \sum_m (u_m(R\mathbf{s}_m))^2$$

- ▶ Linearize $f(R)$ using $\beta_m^{(k)} = u_m(R^{(k)}\mathbf{s}_m)$ and $\delta_m^{(k)} = -\frac{du_m}{dx}(R^{(k)}\mathbf{s}_m) (R^{(k)}\mathbf{s}_m)^\wedge$

$$f(R^{(k+1)}) \approx \frac{1}{2} \sum_m \left(\delta_m^{(k)\top} \boldsymbol{\psi}^{(k)} + \beta_m^{(k)} \right)^2$$

- ▶ The cost is quadratic in $\boldsymbol{\psi}^{(k)}$ and setting its gradient to zero leads to:

$$\left(\sum_m \delta_m^{(k)} \left(\delta_m^{(k)} \right)^\top \right) \boldsymbol{\psi}^{(k)} = - \sum_m \beta_m^{(k)} \delta_m^{(k)}$$

- ▶ Apply the optimal perturbation $\boldsymbol{\psi}^{(k)}$ to the initial guess $R^{(k)}$ according to the left perturbation scheme:

$$R^{(k+1)} = \exp(\hat{\boldsymbol{\psi}}^{(k)}) R^{(k)} \in SO(3)$$

$\mathfrak{so}(3)$ and $SO(3)$ Identities

$$R = \exp(\hat{\theta}) = \sum_{n=0}^{\infty} \frac{1}{n!} \hat{\theta}^n = I + \left(\frac{\sin \|\theta\|}{\|\theta\|} \right) \hat{\theta} + \left(\frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta}^2 \approx I + \hat{\theta}$$

$$R^{-1} = R^T = \exp(-\hat{\theta}) = \sum_{n=0}^{\infty} \frac{1}{n!} (-\hat{\theta})^n \approx I - \hat{\theta}$$

$$\det(R) = 1$$

$$\hat{\theta}^T = -\hat{\theta}$$

$$\text{tr}(R) = 2 \cos \|\theta\| + 1$$

$$\hat{\theta}\theta = 0$$

$$R\theta = \theta$$

$$(A\theta)^\wedge = \hat{\theta}(\text{tr}(A)I - A) - A^T\hat{\theta}, \quad A \in \mathbb{R}^{3 \times 3}$$

$$R\hat{\theta} = \hat{\theta}R$$

$$\hat{\theta}\hat{\phi} = \phi\theta^T - (\theta^T\phi)I, \quad \phi \in \mathbb{R}^3$$

$$(R\mathbf{m})^\wedge = R\hat{\mathbf{m}}R^T, \quad \mathbf{m} \in \mathbb{R}^3$$

$$\hat{\theta}^{2k+1} = (-\theta^T\theta)^k \hat{\theta}$$

$$\exp((R\mathbf{m})^\wedge) = R \exp(\hat{\mathbf{m}}) R^T$$

$$\hat{\theta}\hat{\phi} - \hat{\phi}\hat{\theta} = (\hat{\theta}\phi)^\wedge$$

$SE(3)$ Geometry

Special Euclidean Lie Algebra $\mathfrak{se}(3)$

- ▶ The Lie algebra of $SE(3)$ is the space of twist matrices:

$$\mathfrak{se}(3) := \left\{ \hat{\xi} := \begin{bmatrix} \hat{\theta} & \rho \\ 0 & 0 \end{bmatrix} \mid \xi = \begin{bmatrix} \rho \\ \theta \end{bmatrix} \in \mathbb{R}^6 \right\}$$

- ▶ The **Lie bracket** of $\mathfrak{se}(3)$ is:

$$[\hat{\xi}_1, \hat{\xi}_2] = \hat{\xi}_1 \hat{\xi}_2 - \hat{\xi}_2 \hat{\xi}_1 = \left(\overset{\wedge}{\xi}_1 \xi_2 \right)^\wedge \in \mathfrak{se}(3) \quad \overset{\wedge}{\xi} := \begin{bmatrix} \hat{\theta} & \hat{\rho} \\ 0 & \hat{\theta} \end{bmatrix} \in \mathbb{R}^{6 \times 6}$$

- ▶ The elements $T \in SE(3)$ are related to the elements $\hat{\xi} \in \mathfrak{se}(3)$ through the exponential map:

$$T = \exp(\hat{\xi}) = \sum_{n=0}^{\infty} \frac{1}{n!} (\hat{\xi})^n \quad \xi = \log(T)^\vee$$

Exponential Map from $\mathfrak{se}(3)$ to $SE(3)$

- ▶ The exponential map is **surjective** but **not injective**, i.e., every element of $SE(3)$ can be generated from multiple elements of $\mathfrak{se}(3)$

- ▶ **Rodrigues Formula**: obtained using $\hat{\xi}^4 + \|\theta\|^2 \hat{\xi}^2 = 0$:

$$\begin{aligned} T = \exp(\hat{\xi}) &= \begin{bmatrix} \exp(\hat{\theta}) & J_L(\theta)\rho \\ \mathbf{0}^T & 1 \end{bmatrix} = \sum_{n=0}^{\infty} \frac{1}{n!} \hat{\xi}^n = \\ &= I + \hat{\xi} + \left(\frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\xi}^2 + \left(\frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) \hat{\xi}^3 \end{aligned}$$

- ▶ **Logarithm map** $\log : SE(3) \rightarrow \mathfrak{se}(3)$: for any $T \in SE(3)$, there exists a (not unique) $\xi \in \mathbb{R}^6$ such that:

$$\xi = \begin{bmatrix} \rho \\ \theta \end{bmatrix} = \log(T)^\vee := \begin{cases} \theta = \log(R)^\vee, \rho = J_L^{-1}(\theta)\mathbf{p}, & \text{if } R \neq I, \\ \theta = 0, \rho = \mathbf{p}, & \text{if } R = I. \end{cases}$$

$SE(3)$ Jacobians

- ▶ **Left Jacobian of $SE(3)$:** $\mathcal{J}_L(\xi) = \begin{bmatrix} J_L(\theta) & Q_L(\xi) \\ 0 & J_L(\theta) \end{bmatrix}$
- ▶ **Right Jacobian of $SE(3)$:** $\mathcal{J}_R(\xi) = \begin{bmatrix} J_R(\theta) & Q_R(\xi) \\ 0 & J_R(\theta) \end{bmatrix}$
- ▶ **Baker-Campbell-Hausdorff Formulas:** the $SE(3)$ Jacobians relate small perturbations in $\mathfrak{se}(3)$ to small perturbations in $SE(3)$:

$$\begin{aligned} \exp((\xi + \delta\xi)^\wedge) &\approx \exp(\hat{\xi}) \exp((\mathcal{J}_R(\xi)\delta\xi)^\wedge) \\ &\approx \exp((\mathcal{J}_L(\xi)\delta\xi)^\wedge) \exp(\hat{\xi}) \end{aligned}$$

$$\log(\exp(\hat{\xi}_1) \exp(\hat{\xi}_2))^\vee \approx \begin{cases} \mathcal{J}_L(\xi_2)^{-1}\xi_1 + \xi_2 & \text{if } \xi_1 \text{ is small} \\ \xi_1 + \mathcal{J}_R(\xi_1)^{-1}\xi_2 & \text{if } \xi_2 \text{ is small} \end{cases}$$

Closed-forms of the $SE(3)$ Jacobians

$$\begin{aligned}
 \mathcal{J}_L(\xi) &= \begin{bmatrix} J_L(\theta) & Q_L(\xi) \\ 0 & J_L(\theta) \end{bmatrix} \\
 &= I + \left(\frac{4 - \|\theta\| \sin \|\theta\| - 4 \cos \|\theta\|}{2\|\theta\|^2} \right) \xi^\wedge + \left(\frac{4\|\theta\| - 5 \sin \|\theta\| + \|\theta\| \cos \|\theta\|}{2\|\theta\|^3} \right) \xi^{\wedge^2} \\
 &\quad + \left(\frac{2 - \|\theta\| \sin \|\theta\| - 2 \cos \|\theta\|}{2\|\theta\|^4} \right) \xi^{\wedge^3} + \left(\frac{2\|\theta\| - 3 \sin \|\theta\| + \|\theta\| \cos \|\theta\|}{2\|\theta\|^5} \right) \xi^{\wedge^4} \\
 &\approx I + \frac{1}{2} \xi^\wedge
 \end{aligned}$$

$$\mathcal{J}_L(\xi)^{-1} = \begin{bmatrix} J_L(\theta)^{-1} & -J_L(\theta)^{-1} Q_L(\xi) J_L(\theta)^{-1} \\ \mathbf{0} & J_L(\theta)^{-1} \end{bmatrix} \approx I - \frac{1}{2} \xi^\wedge$$

$$\begin{aligned}
 Q_L(\xi) &= \frac{1}{2} \hat{\rho} + \left(\frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) (\hat{\theta} \hat{\rho} + \hat{\rho} \hat{\theta} + \hat{\theta} \hat{\rho} \hat{\theta}) \\
 &\quad + \left(\frac{\|\theta\|^2 + 2 \cos \|\theta\| - 2}{2\|\theta\|^4} \right) (\hat{\theta}^2 \hat{\rho} + \hat{\rho} \hat{\theta}^2 - 3 \hat{\theta} \hat{\rho} \hat{\theta}) \\
 &\quad + \left(\frac{2\|\theta\| - 3 \sin \|\theta\| + \|\theta\| \cos \|\theta\|}{2\|\theta\|^5} \right) (\hat{\theta} \hat{\rho} \hat{\theta}^2 + \hat{\theta}^2 \hat{\rho} \hat{\theta})
 \end{aligned}$$

$$Q_R(\xi) = Q_L(-\xi) = R Q_L(\xi) + (J_L(\theta) \rho)^\wedge R J_L(\theta)$$

Adjoint

- ▶ The **adjoint** of $T = \begin{bmatrix} R & p \\ \mathbf{0}^T & 1 \end{bmatrix} \in SE(3)$ is:

$$\mathcal{T} = Ad(T) = \begin{bmatrix} R & \hat{p}R \\ \mathbf{0} & R \end{bmatrix} \in \mathbb{R}^{6 \times 6}$$

- ▶ $Ad(SE(3)) := \{\mathcal{T} = Ad(T) \mid T \in SE(3)\}$ is a matrix Lie group

- ▶ The adjoint of $\hat{\xi} = \begin{bmatrix} \hat{\theta} & \rho \\ \mathbf{0}^T & 0 \end{bmatrix} \in \mathfrak{se}(3)$ is:

$$ad(\hat{\xi}) = \overset{\wedge}{\xi} = \begin{bmatrix} \hat{\theta} & \hat{\rho} \\ \mathbf{0} & \hat{\theta} \end{bmatrix} \in \mathbb{R}^{6 \times 6}$$

- ▶ $ad(\mathfrak{se}(3)) := \{\Phi = ad(\hat{\xi}) \in \mathbb{R}^{6 \times 6} \mid \hat{\xi} \in \mathfrak{se}(3)\}$ is the Lie algebra associated with $Ad(SE(3))$

- ▶ The relationship between $\overset{\wedge}{\xi}$ and \mathcal{T} is specified by the exponential map:

$$\mathcal{T} = \exp\left(\overset{\wedge}{\xi}\right) = I + \overset{\wedge}{\xi} \mathcal{J}_L(\xi) \quad \mathcal{J}_L(\xi) = \mathcal{T} \mathcal{J}_R(\xi) = \mathcal{J}_R(-\xi)$$

Pose Lie Groups and Lie Algebras

	Lie algebra		Lie group
4×4	$\xi^\wedge \in \mathfrak{se}(3)$	$\xrightarrow{\text{exp}}$	$\mathbf{T} \in SE(3)$
	$\downarrow \text{ad}$		$\downarrow \text{Ad}$
6×6	$\xi^\wedge \in \text{ad}(\mathfrak{se}(3))$	$\xrightarrow{\text{exp}}$	$\mathcal{T} \in \text{Ad}(SE(3))$

$$\begin{aligned}
 \mathcal{T} &= \underbrace{\text{Ad} \left(\exp(\hat{\xi}) \right)}_{\mathcal{T}} = \exp \left(\underbrace{\text{ad}(\hat{\xi})}_{\xi^\wedge} \right) & \xi &= \begin{bmatrix} \rho \\ \theta \end{bmatrix} \in \mathbb{R}^6 \\
 &= \text{Ad} \left(\exp \left(\begin{bmatrix} \hat{\theta} & \rho \\ \mathbf{0}^T & 0 \end{bmatrix} \right) \right) = \exp \left(\text{ad} \left(\begin{bmatrix} \hat{\theta} & \rho \\ \mathbf{0}^T & 0 \end{bmatrix} \right) \right) \\
 &= \text{Ad} \left(\begin{bmatrix} \exp(\hat{\theta}) & J_L(\theta)\rho \\ \mathbf{0}^T & 1 \end{bmatrix} \right) = \exp \left(\begin{bmatrix} \hat{\theta} & \hat{\rho} \\ \mathbf{0} & \hat{\theta} \end{bmatrix} \right) \\
 &= \begin{bmatrix} \exp(\hat{\theta}) & (J_L(\theta)\rho)^\wedge \exp(\hat{\theta}) \\ \mathbf{0} & \exp(\hat{\theta}) \end{bmatrix}
 \end{aligned}$$

Rodrigues Formula for the Adjoint of $SE(3)$

- ▶ The exponential map is **surjective** but **not injective**, i.e., every element of $Ad(SE(3))$ can be generated from multiple elements of $ad(\mathfrak{se}(3))$
- ▶ **Rodrigues Formula:** using $(\hat{\xi})^5 + 2\|\theta\|^2(\hat{\xi})^3 + \|\theta\|^4\hat{\xi} = 0$ we can obtain a direct expression of $\mathcal{T} \in Ad(SE(3))$ in terms of $\xi = \begin{bmatrix} \rho \\ \theta \end{bmatrix} \in \mathbb{R}^6$:

$$\begin{aligned} \mathcal{T} = Ad(T) &= \exp\left(\begin{matrix} \hat{\xi} \\ \mathbf{0} \end{matrix}\right) = \begin{bmatrix} \exp(\hat{\theta}) & (J_L(\theta)\rho)^\wedge \exp(\hat{\theta}) \\ \mathbf{0} & \exp(\hat{\theta}) \end{bmatrix} = \sum_{n=0}^{\infty} \frac{1}{n!} (\hat{\xi})^n \\ &= I + \left(\frac{3 \sin \|\theta\| - \|\theta\| \cos \|\theta\|}{2\|\theta\|}\right) \hat{\xi} + \left(\frac{4 - \|\theta\| \sin \|\theta\| - 4 \cos \|\theta\|}{2\|\theta\|^2}\right) (\hat{\xi})^2 \\ &\quad + \left(\frac{\sin \|\theta\| - \|\theta\| \cos \|\theta\|}{2\|\theta\|^3}\right) (\hat{\xi})^3 + \left(\frac{2 - \|\theta\| \sin \|\theta\| - 2 \cos \|\theta\|}{2\|\theta\|^4}\right) (\hat{\xi})^4 \end{aligned}$$

Distances in $SE(3)$

- ▶ Two ways to define differences between $SE(3)$ and $Ad(SE(3))$ elements:

$$\xi_{12} = \log (T_1^{-1} T_2)^\vee = \log (T_1^{-1} T_2)^\gamma$$

$$\xi_{21} = \log (T_2 T_1^{-1})^\vee = \log (T_2 T_1^{-1})^\gamma$$

- ▶ Inner product on $\mathfrak{se}(3)$ and $ad(\mathfrak{se}(3))$:

$$\langle \hat{\xi}_1, \hat{\xi}_2 \rangle = -\text{tr} \left(\hat{\xi}_1 \begin{bmatrix} \frac{1}{2} I & \mathbf{0} \\ \mathbf{0}^\top & 1 \end{bmatrix} \hat{\xi}_2^\top \right) = \xi_1^\top \xi_2$$

$$\langle \overset{\wedge}{\xi}_1, \overset{\wedge}{\xi}_2 \rangle = -\text{tr} \left(\overset{\wedge}{\xi}_1 \begin{bmatrix} \frac{1}{4} I & \mathbf{0} \\ \mathbf{0} & \frac{1}{2} I \end{bmatrix} \overset{\wedge}{\xi}_2^\top \right) = \xi_1^\top \xi_2$$

- ▶ The right and left distances on $SE(3)$ and $Ad(SE(3))$ are:

$$\xi_{12} = \sqrt{\langle \hat{\xi}_{12}, \hat{\xi}_{12} \rangle} = \sqrt{\langle \overset{\wedge}{\xi}_{12}, \overset{\wedge}{\xi}_{12} \rangle} = \sqrt{\xi_{12}^\top \xi_{12}} = \|\xi_{12}\|$$

$$\xi_{21} = \sqrt{\langle \hat{\xi}_{21}, \hat{\xi}_{21} \rangle} = \sqrt{\langle \overset{\wedge}{\xi}_{21}, \overset{\wedge}{\xi}_{21} \rangle} = \sqrt{\xi_{21}^\top \xi_{21}} = \|\xi_{21}\|$$

Integration in $SE(3)$

- ▶ The distance between a pose $T = \exp(\hat{\xi})$ and a small perturbation $\exp((\xi + \delta\xi)^\wedge)$ can be approximated using the BCH formulas:

$$\log \left(\exp(\hat{\xi})^{-1} \exp((\xi + \delta\xi)^\wedge) \right)^\vee \approx \mathcal{J}_R(\xi) \delta\xi$$

$$\log \left(\exp((\xi + \delta\xi)^\wedge) \exp(\hat{\xi})^{-1} \right)^\vee \approx \mathcal{J}_L(\xi) \delta\xi$$

- ▶ $|\det(\mathcal{J}(\xi))| = |\det(J(\theta))|^2 = 4 \left(\frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right)^2$
- ▶ Integrating functions of poses can then be carried out as follows:

$$\int_{SE(3)} f(T) dT = \int_{\mathbb{R}^3, \|\theta\| < \pi} f(\xi) |\det(\mathcal{J}(\xi))| d\xi$$

Lie Algebra $\mathfrak{se}(3)$ Identities

$$\hat{\xi} = \begin{bmatrix} \hat{\rho} \\ \theta \end{bmatrix} = \begin{bmatrix} \hat{\theta} & \rho \\ \mathbf{0}^T & 0 \end{bmatrix} \in \mathbb{R}^{4 \times 4} \quad \overset{\wedge}{\xi} = ad(\hat{\xi}) = \begin{bmatrix} \overset{\wedge}{\rho} \\ \theta \end{bmatrix} = \begin{bmatrix} \hat{\theta} & \hat{\rho} \\ \mathbf{0} & \hat{\theta} \end{bmatrix} \in \mathbb{R}^{6 \times 6}$$

$$\overset{\wedge}{\zeta} \overset{\wedge}{\xi} = -\overset{\wedge}{\xi} \overset{\wedge}{\zeta} \quad \zeta \in \mathbb{R}^6$$

$$\overset{\wedge}{\xi} \overset{\wedge}{\xi} = 0$$

$$\hat{\xi}^4 + (\mathbf{m}^T \mathbf{m}) \hat{\xi}^2 = 0 \quad \mathbf{m} \in \mathbb{R}^3$$

$$\left(\overset{\wedge}{\xi}\right)^5 + 2(\mathbf{m}^T \mathbf{m}) \left(\overset{\wedge}{\xi}\right)^3 + (\mathbf{m}^T \mathbf{m})^2 \overset{\wedge}{\xi} = 0$$

$$\mathbf{m}^\odot = \begin{bmatrix} s \\ \lambda \end{bmatrix}^\odot = \begin{bmatrix} \lambda I & -\hat{s} \\ \mathbf{0}^T & \mathbf{0}^T \end{bmatrix} \in \mathbb{R}^{4 \times 6} \quad \mathbf{m}^\odot = \begin{bmatrix} s \\ \lambda \end{bmatrix}^\odot = \begin{bmatrix} 0 & s \\ -\hat{s} & 0 \end{bmatrix} \in \mathbb{R}^{6 \times 4}$$

$$\hat{\xi} \mathbf{m} = \mathbf{m}^\odot \overset{\wedge}{\xi} \quad \mathbf{m}^T \hat{\xi} = \overset{\wedge}{\xi}^T \mathbf{m}^\odot$$

Lie Group $SE(3)$ Identities

$$T = \exp(\hat{\xi}) = \begin{bmatrix} \exp(\hat{\theta}) & J_L(\theta)\rho \\ \mathbf{0}^T & 1 \end{bmatrix}$$
$$= \sum_{n=0}^{\infty} \frac{1}{n!} \hat{\xi}^n = I + \hat{\xi} + \left(\frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\xi}^2 + \left(\frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) \hat{\xi}^3 \approx I + \hat{\xi}$$

$$T^{-1} = \exp(-\hat{\xi}) = \begin{bmatrix} \exp(-\hat{\theta}) & -\exp(-\hat{\theta}) J_L(\theta)\rho \\ \mathbf{0}^T & 1 \end{bmatrix} = \sum_{n=0}^{\infty} \frac{1}{n!} (-\hat{\xi})^n \approx I - \hat{\xi}$$

$$\det(T) = 1$$

$$\text{tr}(T) = 2 \cos \|\theta\| + 2$$

$$T\hat{\xi} = \hat{\xi}T$$

Lie Group $Ad(SE(3))$ Identities

$$\begin{aligned} \mathcal{T} = Ad(T) = \exp \left(\overset{\wedge}{\xi} \right) &= \begin{bmatrix} \exp(\hat{\theta}) & (J_L(\theta)\rho)^\wedge \exp(\hat{\theta}) \\ \mathbf{0} & \exp(\hat{\theta}) \end{bmatrix} \\ &= \sum_{n=0}^{\infty} \frac{1}{n!} \overset{\wedge}{\xi}^n = I + \left(\frac{3 \sin \|\theta\| - \|\theta\| \cos \|\theta\|}{2\|\theta\|} \right) \overset{\wedge}{\xi} + \left(\frac{4 - \|\theta\| \sin \|\theta\| - 4 \cos \|\theta\|}{2\|\theta\|^2} \right) (\overset{\wedge}{\xi})^2 \\ &\quad + \left(\frac{\sin \|\theta\| - \|\theta\| \cos \|\theta\|}{2\|\theta\|^3} \right) (\overset{\wedge}{\xi})^3 + \left(\frac{2 - \|\theta\| \sin \|\theta\| - 2 \cos \|\theta\|}{2\|\theta\|^4} \right) (\overset{\wedge}{\xi})^4 \approx I + \overset{\wedge}{\xi} \end{aligned}$$

$$\mathcal{T}^{-1} = \exp \left(-\overset{\wedge}{\xi} \right) = \begin{bmatrix} \exp(-\hat{\theta}) & -\exp(-\hat{\theta}) (J_L(\theta)\rho)^\wedge \\ \mathbf{0} & \exp(-\hat{\theta}) \end{bmatrix} = \sum_{n=0}^{\infty} \frac{1}{n!} \left(-\overset{\wedge}{\xi} \right)^n \approx I - \overset{\wedge}{\xi}$$

$$\mathcal{T}\xi = \xi$$

$$\mathcal{T}\overset{\wedge}{\xi} = \overset{\wedge}{\xi}\mathcal{T}$$

$$(\mathcal{T}\zeta)^\wedge = \mathcal{T}\hat{\zeta}\mathcal{T}^{-1}$$

$$(\overset{\wedge}{\mathcal{T}\zeta}) = \overset{\wedge}{\mathcal{T}\zeta}\mathcal{T}^{-1} \quad \zeta \in \mathbb{R}^6$$

$$\exp((\mathcal{T}\zeta)^\wedge) = \mathcal{T} \exp(\hat{\zeta}) \mathcal{T}^{-1}$$

$$\exp\left(\overset{\wedge}{(\mathcal{T}\zeta)}\right) = \mathcal{T} \exp\left(\overset{\wedge}{\zeta}\right) \mathcal{T}^{-1}$$

$$(T\mathbf{m})^\odot = T\mathbf{m}^\odot T^{-1}$$

$$((T\mathbf{m})^\odot)^T (T\mathbf{m})^\odot = \mathcal{T}^{-T} (\mathbf{m}^\odot)^T \mathbf{m}^\odot \mathcal{T}^{-1}$$

$SO(3)$ and $SE(3)$ Kinematics

Rotation Kinematics

- ▶ The trajectory $R(t)$ of a continuous rotation motion should satisfy:

$$R^{\top}(t)R(t) = I \quad \Rightarrow \quad \dot{R}^{\top}(t)R(t) + R^{\top}(t)\dot{R}(t) = 0.$$

- ▶ The matrix $R^{\top}(t)\dot{R}(t)$ is **skew-symmetric** and there must exist some vector-valued function $\omega(t) \in \mathbb{R}^3$ such that:

$$R^{\top}(t)\dot{R}(t) = \hat{\omega}(t) \quad \Rightarrow \quad \boxed{\dot{R}(t) = R(t)\hat{\omega}(t)}$$

- ▶ A skew-symmetric matrix gives a first order approximation to a rotation matrix:

$$R(t + dt) \approx R(t) + R(t)\hat{\omega}(t)dt$$

Rotation Kinematics

- ▶ Let $R \in SO(3)$ be the orientation of a rigid body rotating with angular velocity $\omega \in \mathbb{R}^3$ with respect to the world frame.
- ▶ **Rotation kinematic equations of motion:**

$$\dot{R} = R\hat{\omega}_B = \hat{\omega}_W R$$

where ω_B and $\omega_W := R\omega_B$ are the body-frame and world-frame coordinates of ω , respectively.

- ▶ Assuming ω is constant over a short period τ :

$$R(t + \tau) = R(t) \exp(\tau\hat{\omega}_B) = \exp(\tau\hat{\omega}_W)R(t)$$

- ▶ **Discrete Rotation Kinematics:** let $R_k := R(t_k)$, $\tau_k := t_{k+1} - t_k$, and $\omega_k := \omega_B(t_k)$ leading to:

$$R_{k+1} = R_k \exp(\tau_k \hat{\omega}_k)$$

Pose Kinematics

▶ **Angular velocity:** $R^\top(t)\dot{R}(t) = I \Rightarrow R^\top(t)\dot{R}(t) = \hat{\omega}(t) \in \mathfrak{so}(3)$

▶ **Twist:** similarly for $T(t) \in SE(3)$ consider:

$$T^{-1}(t)\dot{T}(t) = \begin{bmatrix} R^\top(t)\dot{R}(t) & R^\top(t)\dot{\mathbf{p}}(t) \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} \hat{\omega}(t) & \mathbf{v}(t) \\ 0 & 0 \end{bmatrix} \in \mathfrak{se}(3)$$

where $\hat{\omega}(t) := R^\top(t)\dot{R}(t)$ and $\mathbf{v}(t) := R^\top(t)\dot{\mathbf{p}}(t)$ are the **body-frame** angular and linear velocities of the body

▶ **Generalized velocity:** $\zeta(t) := \begin{bmatrix} \mathbf{v}(t) \\ \omega(t) \end{bmatrix} \in \mathbb{R}^6$

▶ $\zeta(t)$ is the velocity of the body frame moving relative to the world frame as viewed in the **body frame**

▶ **Continuous-time Pose Kinematics:** $\dot{T}(t) = T(t)\hat{\zeta}(t)$

▶ **Discrete-time Pose Kinematics:** $T_{k+1} = T_k \exp\left(\tau_k \hat{\zeta}_k\right)$

Pose Kinematics

- ▶ Consider a moving body frame $\{B\}$ with pose $T(t) \in SE(3)$
- ▶ Let $\mathbf{s}_B \in \mathbb{R}^3$ be a point in the body frame with homogeneous coordinates $\underline{\mathbf{s}}_B$
- ▶ The velocity of \mathbf{s}_B with respect to the world frame $\{W\}$ can be determined as follows:

$$\begin{aligned}\underline{\mathbf{s}}_W(t) &= T(t)\underline{\mathbf{s}}_B \\ \dot{\underline{\mathbf{s}}}_W(t) &= \dot{T}(t)\underline{\mathbf{s}}_B = \dot{T}(t)T(t)^{-1}\underline{\mathbf{s}}_W(t) \\ &= T(t)\hat{\zeta}(t)T(t)^{-1}\underline{\mathbf{s}}_W(t) \\ &= \begin{bmatrix} R(t)\hat{\omega}(t)R(t)^\top & R(t)\mathbf{v}(t) - R(t)\hat{\omega}(t)R(t)^\top \mathbf{p}(t) \\ \mathbf{0}^\top & 0 \end{bmatrix} \begin{bmatrix} \mathbf{s}_W(t) \\ 1 \end{bmatrix} \\ &= \begin{bmatrix} (R(t)\boldsymbol{\omega}(t))^\wedge (\mathbf{s}_W(t) - \mathbf{p}(t)) + R(t)\mathbf{v}(t) \\ 1 \end{bmatrix}\end{aligned}$$