

# ECE276A: Sensing & Estimation in Robotics

## Lecture 12: $SO(3)$ and $SE(3)$ Geometry and Kinematics

Instructor:

Nikolay Atanasov: [natanasov@ucsd.edu](mailto:natanasov@ucsd.edu)

Teaching Assistants:

Mo Shan: [moshan@eng.ucsd.edu](mailto:moshan@eng.ucsd.edu)

Arash Asgharivaskasi: [aasghari@eng.ucsd.edu](mailto:aasghari@eng.ucsd.edu)

**UC San Diego**

**JACOBS SCHOOL OF ENGINEERING**  
Electrical and Computer Engineering

## Special Orthogonal Group $SO(3)$

- ▶ The orientation  $R$  of a rigid body can be described by a matrix in the **special orthogonal group**:

$$SO(3) := \left\{ R \in \mathbb{R}^{3 \times 3} \mid \underbrace{R^T R = I}_{\text{distances preserved}}, \underbrace{\det(R) = 1}_{\text{no reflection}} \right\}$$

- ▶ It can be verified that  $SO(3)$  satisfies all requirements of a group:
  - ▶ **Closure**:  $R_1 R_2 \in SO(3)$
  - ▶ **Identity**:  $I \in SO(3)$
  - ▶ **Inverse**:  $R^{-1} = R^T \in SO(3)$
  - ▶ **Associativity**:  $(R_1 R_2) R_3 = R_1 (R_2 R_3)$  for all  $R_1, R_2, R_3 \in SO(3)$

# Parametrizations of $SO(3)$

- ▶ **Rotation Matrix:** an element of the **Special Orthogonal Group**:

$$R \in SO(3) := \left\{ R \in \mathbb{R}^{3 \times 3} \mid R^\top R = I, \det(R) = 1 \right\}$$

- ▶ **Euler Angles:** roll  $\phi$ , pitch  $\theta$ , yaw  $\psi$  specifying a **rzyx** rotation:

$$R = R_z(\psi)R_y(\theta)R_x(\phi)$$

- ▶ **Axis-Angle:**  $\theta \in \mathbb{R}^3$  specifying a rotation about an axis  $\eta := \frac{\theta}{\|\theta\|}$  through an angle  $\theta := \|\theta\|$ :

$$R = \exp(\hat{\theta}) = I + \hat{\theta} + \frac{1}{2!}\hat{\theta}^2 + \frac{1}{3!}\hat{\theta}^3 + \dots$$

- ▶ **Unit Quaternion:**  $\mathbf{q} = [q_s, \mathbf{q}_v] \in \{q \in \mathbb{H} \mid q_s^2 + \mathbf{q}_v^\top \mathbf{q}_v = 1\}$ :

$$R = E(\mathbf{q})G(\mathbf{q})^\top \quad \begin{aligned} E(\mathbf{q}) &= [-\mathbf{q}_v, q_s I + \hat{\mathbf{q}}_v] \\ G(\mathbf{q}) &= [-\mathbf{q}_v, q_s I - \hat{\mathbf{q}}_v] \end{aligned}$$

# Special Euclidean Group $SE(3)$

- ▶ The pose  $T$  of a rigid body can be described by a matrix in the **special Euclidean group**:

$$SE(3) := \left\{ T := \begin{bmatrix} R & \mathbf{p} \\ \mathbf{0}^\top & 1 \end{bmatrix} \in \mathbb{R}^{4 \times 4} \mid R \in SO(3), \mathbf{p} \in \mathbb{R}^3 \right\}$$

- ▶ It can be verified that  $SE(3)$  satisfies all requirements of a group:
  - ▶ **Closure:**  $T_1 T_2 = \begin{bmatrix} R_1 & \mathbf{p}_1 \\ \mathbf{0}^\top & 1 \end{bmatrix} \begin{bmatrix} R_2 & \mathbf{p}_2 \\ \mathbf{0}^\top & 1 \end{bmatrix} = \begin{bmatrix} R_1 R_2 & R_1 \mathbf{p}_2 + \mathbf{p}_1 \\ \mathbf{0}^\top & 1 \end{bmatrix} \in SE(3)$
  - ▶ **Identity:**  $\begin{bmatrix} I & \mathbf{0} \\ \mathbf{0}^\top & 1 \end{bmatrix} \in SE(3)$
  - ▶ **Inverse:**  $\begin{bmatrix} R & \mathbf{p} \\ \mathbf{0}^\top & 1 \end{bmatrix}^{-1} = \begin{bmatrix} R^\top & -R^\top \mathbf{p} \\ \mathbf{0}^\top & 1 \end{bmatrix} \in SE(3)$
  - ▶ **Associativity:**  $(T_1 T_2) T_3 = T_1 (T_2 T_3)$  for all  $T_1, T_2, T_3 \in SE(3)$

# Matrix Lie Group

- ▶  $SO(3)$  and  $SE(3)$  are **matrix Lie groups**
- ▶ A **group** is a set of elements with an operation that combines any two elements to form a third element also in the set. A group satisfies four axioms: closure, associativity, identity, and invertibility
- ▶ A **manifold** is a topological space that is locally homeomorphic to Euclidean space but globally may have more complicated structure
- ▶ A **Lie group** is a group that is also a differentiable manifold with the property that the group operations are smooth
- ▶ A **matrix Lie group** further specifies that the group elements are matrices, the combination operation is matrix multiplication, and the inversion operation is matrix inversion
- ▶ The **exponential map** relates a matrix Lie group to its **Lie algebra**

$$\exp(A) = \sum_{n=0}^{\infty} \frac{1}{n!} A^n \quad \log(A) = \sum_{n=1}^{\infty} \frac{(-1)^{n-1}}{n} (A - I)^n$$

# Lie Algebra

- ▶ A **Lie algebra** is a vector space  $\mathbb{V}$  over some field  $\mathbb{F}$  with a binary operation,  $[\cdot, \cdot]$ , called a **Lie bracket**
- ▶ For all  $X, Y, Z \in \mathbb{V}$  and  $a, b \in \mathbb{F}$ , the Lie bracket satisfies:

$$\text{closure :} \quad [X, Y] \in \mathbb{V}$$

$$\text{bilinearity :} \quad [aX + bY, Z] = a[X, Z] + b[Y, Z]$$

$$[Z, aX + bY] = a[Z, X] + b[Z, Y]$$

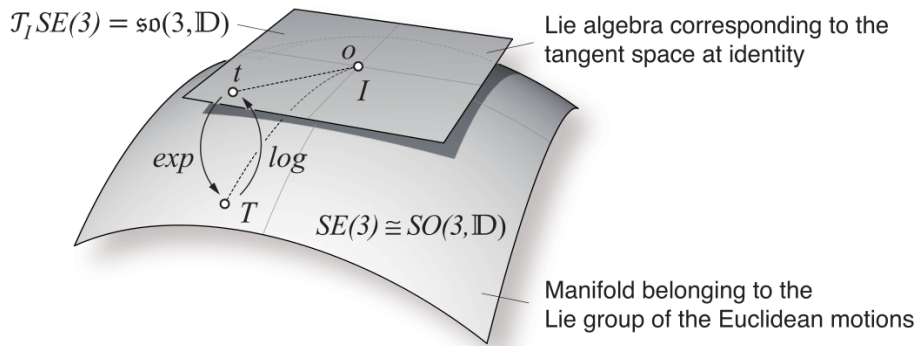
$$\text{alternating :} \quad [X, X] = 0$$

$$\text{Jacobi identity :} \quad [X, [Y, Z]] + [Y, [Z, X]] + [Z, [X, Y]] = 0$$

- ▶ A **Lie algebra** may be associated with every Lie group. The vector space of a Lie algebra forms the **tangent space** to the Lie group at the identity element of the group.

# Lie Group and Lie Algebra Visualization

- **Lie Group:** free of singularities but has constraints
- **Lie Algebra:** free of constraints but has singularities



**Figure:**  $SE(3)$  and the corresponding Lie algebra  $\mathfrak{se}(3)$  as tangent space at identity

## $SO(3)$ Geometry



# Special Orthogonal Lie Algebra $\mathfrak{so}(3)$

- ▶ The **Lie algebra** of  $SO(3)$  is the space of skew-symmetric matrices

$$\mathfrak{so}(3) := \{\hat{\boldsymbol{\theta}} \in \mathbb{R}^{3 \times 3} \mid \boldsymbol{\theta} \in \mathbb{R}^3\}$$

- ▶ The **Lie bracket** of  $\mathfrak{so}(3)$  is:

$$[\hat{\boldsymbol{\theta}}_1, \hat{\boldsymbol{\theta}}_2] = \hat{\boldsymbol{\theta}}_1 \hat{\boldsymbol{\theta}}_2 - \hat{\boldsymbol{\theta}}_2 \hat{\boldsymbol{\theta}}_1 = \left( \hat{\boldsymbol{\theta}}_1 \boldsymbol{\theta}_2 \right)^\wedge \in \mathfrak{so}(3)$$

- ▶ **Generators of  $\mathfrak{so}(3)$ :** derivatives of rotations around each standard axis:

$$G_x = \left. \frac{d}{d\phi} R_x(\phi) \right|_{\phi=0} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{bmatrix} \quad G_y = \begin{bmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{bmatrix} \quad G_z = \begin{bmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}$$

- ▶ The elements  $\hat{\boldsymbol{\theta}} = \theta_1 G_x + \theta_2 G_y + \theta_3 G_z \in \mathfrak{so}(3)$  are linear combinations of the generators

## Exponential Map from $\mathfrak{so}(3)$ to $SO(3)$

- ▶ The elements  $R \in SO(3)$  are related to the elements  $\hat{\theta} \in \mathfrak{so}(3)$  through the **exponential map**:

$$R = \exp(\hat{\theta}) = \sum_{n=0}^{\infty} \frac{1}{n!} (\hat{\theta})^n$$

- ▶ The exponential map is **surjective** but **not injective**, i.e., every element of  $SO(3)$  can be generated from multiple elements of  $\mathfrak{so}(3)$
- ▶ Any vector  $(\|\theta\| + 2\pi k) \frac{\theta}{\|\theta\|}$  for integer  $k$  leads to the same  $R \in SO(3)$
- ▶ The exponential map is **not commutative**,  $e^{\hat{\theta}_1} e^{\hat{\theta}_2} \neq e^{\hat{\theta}_2} e^{\hat{\theta}_1} \neq e^{\hat{\theta}_1 + \hat{\theta}_2}$ , unless  $[\hat{\theta}_1, \hat{\theta}_2] = \hat{\theta}_1 \hat{\theta}_2 - \hat{\theta}_2 \hat{\theta}_1 = 0$

## Rodrigues Formula

- ▶ A closed-form expression for the exponential map from  $\mathfrak{so}(3)$  to  $SO(3)$ :

$$R = \exp(\hat{\theta}) = I + \left( \frac{\sin \|\theta\|}{\|\theta\|} \right) \hat{\theta} + \left( \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta}^2$$

- ▶ The formula is derived using that  $\hat{\theta}^{2n+1} = (-\theta^\top \theta)^n \hat{\theta}$ :

$$\begin{aligned} \exp(\hat{\theta}) &= I + \sum_{n=1}^{\infty} \frac{1}{n!} \hat{\theta}^n \\ &= I + \sum_{n=0}^{\infty} \frac{1}{(2n+1)!} \hat{\theta}^{2n+1} + \sum_{n=0}^{\infty} \frac{1}{(2n+2)!} \hat{\theta}^{2n+2} \\ &= I + \left( \sum_{n=0}^{\infty} \frac{(-1)^n \|\theta\|^{2n}}{(2n+1)!} \right) \hat{\theta} + \left( \sum_{n=0}^{\infty} \frac{(-1)^n \|\theta\|^{2n}}{(2n+2)!} \right) \hat{\theta}^2 \\ &= I + \left( \frac{\sin \|\theta\|}{\|\theta\|} \right) \hat{\theta} + \left( \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta}^2 \end{aligned}$$

## Logarithm Map from $SO(3)$ to $\mathfrak{so}(3)$

- ▶  $\forall R \in SO(3)$ , there exists a (non-unique)  $\theta \in \mathbb{R}^3$  such that  $R = \exp(\hat{\theta})$
- ▶ The **logarithm map**  $\log : SO(3) \rightarrow \mathfrak{so}(3)$  is the inverse of  $\exp(\hat{\theta})$ :

$$\theta = \|\theta\| = \arccos\left(\frac{\text{tr}(R) - 1}{2}\right)$$

$$\eta = \frac{\theta}{\|\theta\|} = \frac{1}{2 \sin(\|\theta\|)} \begin{bmatrix} R_{32} - R_{23} \\ R_{13} - R_{31} \\ R_{21} - R_{12} \end{bmatrix}$$

$$\hat{\theta} = \log(R) = \frac{\|\theta\|}{2 \sin \|\theta\|} (R - R^\top)$$

- ▶ If  $R = I$ , then  $\theta = 0$  and  $\eta$  is undefined

- ▶ If  $\text{tr}(R) = -1$ , then  $\theta = \pi$  and for any  $i \in \{1, 2, 3\}$ :

$$\eta = \frac{1}{\sqrt{2(1 + e_i^\top R e_i)}} (I + R) e_i$$

- ▶ The log map has a singularity at  $\theta = 0$  because there are infinite choices of rotation axes or equivalently the exponential map is many-to-one.
- ▶ The matrix exponential “integrates”  $\hat{\theta} \in \mathfrak{se}(3)$  for one second; the matrix logarithm “differentiates”  $R \in SO(3)$  to obtain  $\hat{\theta} \in \mathfrak{se}(3)$

## $SO(3)$ Jacobians

- ▶ The **left Jacobian** of  $SO(3)$  is the matrix:

$$J_L(\theta) := \sum_{n=0}^{\infty} \frac{1}{(n+1)!} (\hat{\theta})^n \quad R = I + \hat{\theta} J_L(\theta)$$

- ▶ The **right Jacobian** of  $SO(3)$  is the matrix:

$$J_R(\theta) := \sum_{n=0}^{\infty} \frac{1}{(n+1)!} (-\hat{\theta})^n \quad J_R(\theta) = J_L(-\theta) = J_L(\theta)^\top = R^\top J_L(\theta)$$

- ▶ **Baker-Campbell-Hausdorff Formulas:** the  $SO(3)$  Jacobians relate small perturbations in  $\mathfrak{so}(3)$  to small perturbations in  $SO(3)$ :

$$\begin{aligned} \exp((\theta + \delta\theta)^\wedge) &\approx \exp(\hat{\theta}) \exp((J_R(\theta)\delta\theta)^\wedge) \\ &\approx \exp((J_L(\theta)\delta\theta)^\wedge) \exp(\hat{\theta}) \end{aligned}$$

$$\log(\exp(\hat{\theta}_1) \exp(\hat{\theta}_2))^\vee \approx \begin{cases} J_L(\theta_2)^{-1} \theta_1 + \theta_2 & \text{if } \theta_1 \text{ is small} \\ \theta_1 + J_R(\theta_1)^{-1} \theta_2 & \text{if } \theta_2 \text{ is small} \end{cases}$$

## Closed-forms of the $SO(3)$ Jacobians

$$J_L(\theta) = I + \left( \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta} + \left( \frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) \hat{\theta}^2 \approx I + \frac{1}{2} \hat{\theta}$$

$$J_L(\theta)^{-1} = I - \frac{1}{2} \hat{\theta} + \left( \frac{1}{\|\theta\|^2} - \frac{1 + \cos \|\theta\|}{2\|\theta\| \sin \|\theta\|} \right) \hat{\theta}^2 \approx I - \frac{1}{2} \hat{\theta}$$

$$J_R(\theta) = I - \left( \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta} + \left( \frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) \hat{\theta}^2 \approx I - \frac{1}{2} \hat{\theta}$$

$$J_R(\theta)^{-1} = I + \frac{1}{2} \hat{\theta} + \left( \frac{1}{\|\theta\|^2} - \frac{1 + \cos \|\theta\|}{2\|\theta\| \sin \|\theta\|} \right) \hat{\theta}^2 \approx I + \frac{1}{2} \hat{\theta}$$

$$J_L(\theta) J_L(\theta)^T = I + \left( 1 - 2 \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta}^2 \succ 0$$

$$\left( J_L(\theta) J_L(\theta)^T \right)^{-1} = I + \left( 1 - 2 \frac{\|\theta\|^2}{1 - \cos \|\theta\|} \right) \hat{\theta}^2$$

## Distances in $SO(3)$

- There are two ways to define the difference between two rotations:

$$\boldsymbol{\theta}_{12} = \log \left( R_1^\top R_2 \right)^\vee \quad \boldsymbol{\theta}_{21} = \log \left( R_2 R_1^\top \right)^\vee \quad R_1, R_2 \in SO(3)$$

- Inner product on  $\mathfrak{so}(3)$ :

$$\langle \hat{\boldsymbol{\theta}}_1, \hat{\boldsymbol{\theta}}_2 \rangle = \frac{1}{2} \operatorname{tr} \left( \hat{\boldsymbol{\theta}}_1^\top \hat{\boldsymbol{\theta}}_2 \right) = \boldsymbol{\theta}_1^\top \boldsymbol{\theta}_2$$

- The metric distance between two rotations may be defined in two ways as the magnitude of the rotation difference:

$$\theta_{12} := \sqrt{\langle \log \left( R_1^\top R_2 \right), \log \left( R_1^\top R_2 \right) \rangle} = \|\boldsymbol{\theta}_{12}\|_2$$

$$\theta_{21} := \sqrt{\langle \log \left( R_2 R_1^\top \right), \log \left( R_2 R_1^\top \right) \rangle} = \|\boldsymbol{\theta}_{21}\|_2$$

## Integration in $SO(3)$

- ▶ The distance between a rotation  $R = \exp(\hat{\theta})$  and a small perturbation  $\exp((\theta + \delta\theta)^\wedge)$  can be approximated using the BCH formulas:

$$\begin{aligned}\log \left( \exp(\hat{\theta})^\top \exp((\theta + \delta\theta)^\wedge) \right)^\vee &\approx \log \left( R^\top R \exp((J_R(\theta)\delta\theta)^\wedge) \right)^\vee = J_R(\theta)\delta\theta \\ \log \left( \exp((\theta + \delta\theta)^\wedge) \exp(\hat{\theta})^\top \right)^\vee &\approx \log \left( \exp((J_L(\theta)\delta\theta)^\wedge) R R^\top \right)^\vee = J_L(\theta)\delta\theta\end{aligned}$$

- ▶ Regardless of which distance metric we use, the infinitesimal volume element is the same:

$$\det(J_L(\theta)) = \det(J_R(\theta)) \quad dR = |\det(J(\theta))| d\theta = 2 \left( \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) d\theta$$

- ▶ Integrating functions of rotations can then be carried out as follows:

$$\int_{SO(3)} f(R) dR = \int_{\|\theta\| < \pi} f(\exp(\hat{\theta})) |\det(J(\theta))| d\theta$$



## Derivatives in $SO(3)$

- Consider  $\mathbf{s} \in \mathbb{R}^3$  rotated by a rotation matrix  $R \in SO(3)$  to a new frame
- How do we compute the derivative of  $R\mathbf{s}$  with respect to the rotation  $R$ ?
- Let  $\boldsymbol{\theta} \in \mathbb{R}^3$  be the Lie algebra vector representing  $R$ , i.e.,  $R = \exp(\hat{\boldsymbol{\theta}})$
- We can compute derivatives with respect to the elements of  $\boldsymbol{\theta}$ :

$$\begin{aligned}
 \frac{\partial R\mathbf{s}}{\partial \theta_i} &= \lim_{h \rightarrow 0} \frac{\exp((\boldsymbol{\theta} + h\mathbf{e}_i)^\wedge) \mathbf{s} - \exp(\hat{\boldsymbol{\theta}}) \mathbf{s}}{h} \\
 &\stackrel{\text{BCH}}{\stackrel{\text{Formula}}{=}} \lim_{h \rightarrow 0} \frac{\exp((hJ_L(\boldsymbol{\theta})\mathbf{e}_i)^\wedge) \exp(\hat{\boldsymbol{\theta}}) \mathbf{s} - \exp(\hat{\boldsymbol{\theta}}) \mathbf{s}}{h} \\
 &\stackrel{\exp(\delta\hat{\boldsymbol{\theta}}) \approx I + \delta\hat{\boldsymbol{\theta}}}{=} \lim_{h \rightarrow 0} \frac{(I + h(J_L(\boldsymbol{\theta})\mathbf{e}_i)^\wedge) \exp(\hat{\boldsymbol{\theta}}) \mathbf{s} - \exp(\hat{\boldsymbol{\theta}}) \mathbf{s}}{h} \\
 &= (J_L(\boldsymbol{\theta})\mathbf{e}_i)^\wedge R\mathbf{s} = -(R\mathbf{s})^\wedge J_L(\boldsymbol{\theta})\mathbf{e}_i
 \end{aligned}$$

- Stacking the three directional derivatives:  $\boxed{\frac{\partial R\mathbf{s}}{\partial \boldsymbol{\theta}} = -(R\mathbf{s})^\wedge J_L(\boldsymbol{\theta})}$

## Derivatives in $SO(3)$

- **Perturbation in  $\mathfrak{so}(3)$ :** the gradient can also be obtained via a small perturbation  $\delta\boldsymbol{\theta}$  to the axis-angle vector  $\boldsymbol{\theta}$ :

$$\begin{aligned}\exp((\boldsymbol{\theta} + \delta\boldsymbol{\theta})^\wedge) \mathbf{s} &\stackrel{\text{BCH}}{\approx} \exp((J_L \boldsymbol{\theta}) \delta\boldsymbol{\theta})^\wedge \exp(\hat{\boldsymbol{\theta}}) \mathbf{s} \\ &\approx (I + (J_L(\boldsymbol{\theta}) \delta\boldsymbol{\theta})^\wedge) \exp(\hat{\boldsymbol{\theta}}) \mathbf{s} \\ &= R\mathbf{s} + (J_L(\boldsymbol{\theta}) \delta\boldsymbol{\theta})^\wedge R\mathbf{s} = R\mathbf{s} - \underbrace{(R\mathbf{s})^\wedge J_L(\boldsymbol{\theta}) \delta\boldsymbol{\theta}}_{\frac{\partial R\mathbf{s}}{\partial \boldsymbol{\theta}}}\end{aligned}$$

- This is the same as using first-order Taylor series to identify the Jacobian of a function  $f(\mathbf{x})$ :

$$f(\mathbf{x} + \delta\mathbf{x}) \approx f(\mathbf{x}) + \left[ \frac{\partial f}{\partial \mathbf{x}}(\mathbf{x}) \right] \delta\mathbf{x}$$

- **Perturbation in  $SO(3)$ :** a small perturbation  $\boldsymbol{\psi} = J_L(\boldsymbol{\theta}) \delta\boldsymbol{\theta}$  may also be applied directly to  $R$ :

$$\exp(\hat{\boldsymbol{\psi}}) R\mathbf{s} \approx (I + \hat{\boldsymbol{\psi}}) R\mathbf{s} = R\mathbf{s} - (R\mathbf{s})^\wedge \boldsymbol{\psi}$$

## Gradient Descent in $SO(3)$

- ▶ Consider  $\min_{\mathbf{x}} f(\mathbf{x})$
- ▶ **Gradient descent in  $\mathbb{R}^d$ :** given an initial guess  $\mathbf{x}^{(k)}$  take a step of size  $\alpha^{(k)} > 0$  along the descent direction  $\delta \mathbf{x}^{(k)} = -\nabla f(\mathbf{x}^{(k)})$ :

$$\mathbf{x}^{(k+1)} = \mathbf{x}^{(k)} + \alpha^{(k)} \delta \mathbf{x}^{(k)}$$

- ▶ Consider  $\min_R f(R\mathbf{s})$
- ▶ **Gradient descent in  $SO(3)$ :** given an initial guess  $R^{(k)}$  take a step of size  $\alpha^{(k)} > 0$  along the descent direction  $\psi^{(k)} = -\delta^{(k)}$ :

$$R^{(k+1)} = \exp\left(\alpha^{(k)} \hat{\psi}^{(k)}\right) R^{(k)}$$

where  $\delta^{(k)}$  should be the gradient of  $f$  wrt  $R$  evaluated at  $R^{(k)}\mathbf{s}$

## Choosing a Descent Direction in $SO(3)$

- Use a perturbation  $\psi^{(k)}$  around the initial guess  $R^{(k)}$  to determine the gradient  $\delta^{(k)}$ :

$$\begin{aligned} f\left(\exp(\hat{\psi}^{(k)})R^{(k)}\mathbf{s}\right) &\approx f\left((I + \hat{\psi}^{(k)})R^{(k)}\mathbf{s}\right) \\ &= f\left(R^{(k)}\mathbf{s} - \left(R^{(k)}\mathbf{s}\right)^{\wedge} \psi^{(k)}\right) \\ &\approx f\left(R^{(k)}\mathbf{s}\right) - \underbrace{\nabla f\left(R^{(k)}\mathbf{s}\right)^{\top} \left(R^{(k)}\mathbf{s}\right)^{\wedge}}_{\delta^{(k)\top}} \psi^{(k)} \end{aligned}$$

- **Gradient descent in  $SO(3)$ :** given an initial guess  $R^{(k)}$  take a step of size  $\alpha^{(k)} > 0$  along the descent direction  $\psi^{(k)} = -\delta^{(k)}$ :

$$\begin{aligned} \psi^{(k)} &= -\left(R^{(k)}\mathbf{s}\right)^{\wedge} \nabla f(R^{(k)}\mathbf{s}) \\ R^{(k+1)} &= \exp\left(\alpha^{(k)}\hat{\psi}^{(k)}\right) R^{(k)} \end{aligned}$$

## Gauss-Newton Optimization in $SO(3)$

- Optimization problem:

$$\min_R f(R) := \frac{1}{2} \sum_j \mathbf{e}_j(R\mathbf{s}_j)^\top \mathbf{e}_j(R\mathbf{s}_j)$$

- Linearize  $f(R)$  using  $\mathbf{e}_j^{(k)} := \mathbf{e}_j(R^{(k)}\mathbf{s}_j)$  and  $J_j^{(k)} := -\frac{d\mathbf{e}_j}{d\mathbf{x}}(R^{(k)}\mathbf{s}_j) (R^{(k)}\mathbf{s}_j)^\wedge$   
 $f(R^{(k+1)}) = f(\exp(\hat{\psi}^{(k)})R^{(k)}) \approx \frac{1}{2} \sum_j \left( \mathbf{e}_j^{(k)} + J_j^{(k)}\psi^{(k)} \right)^\top \left( \mathbf{e}_j^{(k)} + J_j^{(k)}\psi^{(k)} \right)$

- The cost is quadratic in  $\psi^{(k)}$  and setting its gradient to zero leads to:

$$\left( \sum_j J_j^{(k)} (J_j^{(k)})^\top \right) \psi^{(k)} = - \sum_j (J_j^{(k)})^\top \mathbf{e}_j^{(k)}$$

- Apply the optimal perturbation  $\psi^{(k)}$  to the initial guess  $R^{(k)}$  according to the left perturbation scheme:

$$R^{(k+1)} = \exp(\alpha^{(k)} \hat{\psi}^{(k)}) R^{(k)}$$

## $SO(3)$ and $\mathfrak{so}(3)$ Identities

$$R = \exp(\hat{\theta}) = \sum_{n=0}^{\infty} \frac{1}{n!} \hat{\theta}^n = I + \left( \frac{\sin \|\theta\|}{\|\theta\|} \right) \hat{\theta} + \left( \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\theta}^2 \approx I + \hat{\theta}$$

$$R^{-1} = R^{\top} = \exp(-\hat{\theta}) = \sum_{n=0}^{\infty} \frac{1}{n!} (-\hat{\theta})^n \approx I - \hat{\theta}$$

$$\det(R) = 1 \qquad \hat{\theta}^{\top} = -\hat{\theta}$$

$$\text{tr}(R) = 2 \cos \|\theta\| + 1 \qquad \hat{\theta}\theta = 0$$

$$R\theta = \theta \qquad (A\theta)^{\wedge} = \hat{\theta}(\text{tr}(A)I - A) - A^{\top}\hat{\theta}, \quad A \in \mathbb{R}^{3 \times 3}$$

$$R\hat{\theta} = \hat{\theta}R \qquad \hat{\theta}\hat{\phi} = \phi\theta^{\top} - (\theta^{\top}\phi)I, \quad \phi \in \mathbb{R}^3$$

$$(R\mathbf{s})^{\wedge} = R\hat{\mathbf{s}}R^{\top}, \quad \mathbf{s} \in \mathbb{R}^3 \qquad \hat{\theta}^{2k+1} = (-\theta^{\top}\theta)^k \hat{\theta}$$

$$\exp((R\mathbf{s})^{\wedge}) = R \exp(\hat{\mathbf{s}}) R^{\top} \qquad [\theta, \phi] = \hat{\theta}\hat{\phi} - \hat{\phi}\hat{\theta} = (\hat{\theta}\phi)^{\wedge}$$

## $SE(3)$ Geometry

## Special Euclidean Lie Algebra $\mathfrak{se}(3)$

- ▶ The Lie algebra of  $SE(3)$  is the space of twist matrices:

$$\mathfrak{se}(3) := \left\{ \hat{\xi} := \begin{bmatrix} \hat{\theta} & \rho \\ 0 & 0 \end{bmatrix} \in \mathbb{R}^{4 \times 4} \mid \xi = \begin{bmatrix} \rho \\ \theta \end{bmatrix} \in \mathbb{R}^6 \right\}$$

- ▶ The **Lie bracket** of  $\mathfrak{se}(3)$  is:

$$[\hat{\xi}_1, \hat{\xi}_2] = \hat{\xi}_1 \hat{\xi}_2 - \hat{\xi}_2 \hat{\xi}_1 = \left( \overset{\wedge}{\xi}_1 \xi_2 \right)^\wedge \in \mathfrak{se}(3) \quad \overset{\wedge}{\xi} := \begin{bmatrix} \hat{\theta} & \hat{\rho} \\ 0 & \hat{\theta} \end{bmatrix} \in \mathbb{R}^{6 \times 6}$$

- ▶ The elements  $T \in SE(3)$  are related to the elements  $\hat{\xi} \in \mathfrak{se}(3)$  through the exponential map:

$$T = \exp(\hat{\xi}) = \sum_{n=0}^{\infty} \frac{1}{n!} (\hat{\xi})^n \quad \xi = \log(T)^\vee$$



## Exponential Map from $\mathfrak{se}(3)$ to $SE(3)$

- ▶ The exponential map is **surjective** but **not injective**, i.e., every element of  $SE(3)$  can be generated from multiple elements of  $\mathfrak{se}(3)$
- ▶ **Rodrigues Formula**: obtained using  $\hat{\xi}^4 + \|\theta\|^2 \hat{\xi}^2 = 0$ :

$$\begin{aligned} T = \exp(\hat{\xi}) &= \begin{bmatrix} \exp(\hat{\theta}) & J_L(\theta)\rho \\ \mathbf{0}^T & 1 \end{bmatrix} = \sum_{n=0}^{\infty} \frac{1}{n!} \hat{\xi}^n = \\ &= I + \hat{\xi} + \left( \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\xi}^2 + \left( \frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) \hat{\xi}^3 \end{aligned}$$

- ▶ **Logarithm map**  $\log : SE(3) \rightarrow \mathfrak{se}(3)$ : for any  $T \in SE(3)$ , there exists a (non-unique)  $\xi \in \mathbb{R}^6$  such that:

$$\xi = \begin{bmatrix} \rho \\ \theta \end{bmatrix} = \log(T)^\vee := \begin{cases} \theta = \log(R)^\vee, \rho = J_L^{-1}(\theta)\mathbf{p}, & \text{if } R \neq I, \\ \theta = 0, \rho = \mathbf{p}, & \text{if } R = I. \end{cases}$$

## $SE(3)$ Jacobians

- ▶ **Left Jacobian of  $SE(3)$ :**  $\mathcal{J}_L(\xi) = \begin{bmatrix} J_L(\theta) & Q_L(\xi) \\ 0 & J_L(\theta) \end{bmatrix}$
- ▶ **Right Jacobian of  $SE(3)$ :**  $\mathcal{J}_R(\xi) = \begin{bmatrix} J_R(\theta) & Q_R(\xi) \\ 0 & J_R(\theta) \end{bmatrix}$
- ▶ **Baker-Campbell-Hausdorff Formulas:** the  $SE(3)$  Jacobians relate small perturbations in  $\mathfrak{se}(3)$  to small perturbations in  $SE(3)$ :

$$\begin{aligned} \exp((\xi + \delta\xi)^\wedge) &\approx \exp(\hat{\xi}) \exp((\mathcal{J}_R(\xi)\delta\xi)^\wedge) \\ &\approx \exp((\mathcal{J}_L(\xi)\delta\xi)^\wedge) \exp(\hat{\xi}) \end{aligned}$$

$$\log(\exp(\hat{\xi}_1) \exp(\hat{\xi}_2))^\vee \approx \begin{cases} \mathcal{J}_L(\xi_2)^{-1}\xi_1 + \xi_2 & \text{if } \xi_1 \text{ is small} \\ \xi_1 + \mathcal{J}_R(\xi_1)^{-1}\xi_2 & \text{if } \xi_2 \text{ is small} \end{cases}$$

## Closed-forms of the $SE(3)$ Jacobians

$$\begin{aligned}
 \mathcal{J}_L(\xi) &= \sum_{n=0}^{\infty} \frac{1}{(n+1)!} (\overset{\wedge}{\xi})^n = \begin{bmatrix} J_L(\theta) & Q_L(\xi) \\ 0 & J_L(\theta) \end{bmatrix} \\
 &= I + \left( \frac{4 - \|\theta\| \sin \|\theta\| - 4 \cos \|\theta\|}{2\|\theta\|^2} \right) \overset{\wedge}{\xi} + \left( \frac{4\|\theta\| - 5 \sin \|\theta\| + \|\theta\| \cos \|\theta\|}{2\|\theta\|^3} \right) \overset{\wedge}{\xi}^2 \\
 &\quad + \left( \frac{2 - \|\theta\| \sin \|\theta\| - 2 \cos \|\theta\|}{2\|\theta\|^4} \right) \overset{\wedge}{\xi}^3 + \left( \frac{2\|\theta\| - 3 \sin \|\theta\| + \|\theta\| \cos \|\theta\|}{2\|\theta\|^5} \right) \overset{\wedge}{\xi}^4 \\
 &\approx I + \frac{1}{2} \overset{\wedge}{\xi}
 \end{aligned}$$

$$\mathcal{J}_L(\xi)^{-1} = \begin{bmatrix} J_L(\theta)^{-1} & -J_L(\theta)^{-1} Q_L(\xi) J_L(\theta)^{-1} \\ \mathbf{0} & J_L(\theta)^{-1} \end{bmatrix} \approx I - \frac{1}{2} \overset{\wedge}{\xi}$$

$$\begin{aligned}
 Q_L(\xi) &= \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} \frac{1}{(n+m+2)!} \hat{\theta}^n \hat{\rho} \hat{\theta}^m \\
 &= \frac{1}{2} \hat{\rho} + \left( \frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) (\hat{\theta} \hat{\rho} + \hat{\rho} \hat{\theta} + \hat{\theta} \hat{\rho} \hat{\theta}) + \left( \frac{\|\theta\|^2 + 2 \cos \|\theta\| - 2}{2\|\theta\|^4} \right) (\hat{\theta}^2 \hat{\rho} + \hat{\rho} \hat{\theta}^2 - 3 \hat{\theta} \hat{\rho} \hat{\theta}) \\
 &\quad + \left( \frac{2\|\theta\| - 3 \sin \|\theta\| + \|\theta\| \cos \|\theta\|}{2\|\theta\|^5} \right) (\hat{\theta} \hat{\rho} \hat{\theta}^2 + \hat{\theta}^2 \hat{\rho} \hat{\theta})
 \end{aligned}$$

$$Q_R(\xi) = Q_L(-\xi) = R Q_L(\xi) + (J_L(\theta) \rho)^\wedge R J_L(\theta)$$

## Adjoint

- ▶ The **adjoint** of  $T = \begin{bmatrix} R & \mathbf{p} \\ \mathbf{0}^\top & 1 \end{bmatrix} \in SE(3)$  is:

$$Ad(T) := \begin{bmatrix} R & \hat{\mathbf{p}}R \\ \mathbf{0} & R \end{bmatrix} \in \mathbb{R}^{6 \times 6}$$

- ▶  $Ad(SE(3)) := \{Ad(T) \in \mathbb{R}^{6 \times 6} \mid T \in SE(3)\}$  is a matrix Lie group

- ▶ The adjoint of  $\hat{\xi} = \begin{bmatrix} \hat{\theta} & \rho \\ \mathbf{0}^\top & 0 \end{bmatrix} \in \mathfrak{se}(3)$  is:

$$ad(\hat{\xi}) := \overset{\wedge}{\xi} = \begin{bmatrix} \hat{\theta} & \hat{\rho} \\ \mathbf{0} & \hat{\theta} \end{bmatrix} \in \mathbb{R}^{6 \times 6}$$

- ▶  $ad(\mathfrak{se}(3)) := \{ad(\hat{\xi}) \in \mathbb{R}^{6 \times 6} \mid \hat{\xi} \in \mathfrak{se}(3)\}$  is the Lie algebra associated with  $Ad(SE(3))$

- ▶ The relationship between  $\overset{\wedge}{\xi}$  and  $\mathcal{T} := Ad(T)$  is specified by the exponential map:

$$\mathcal{T} = \exp(\overset{\wedge}{\xi}) = I + \overset{\wedge}{\xi} \mathcal{J}_L(\xi) \quad \mathcal{J}_L(\xi) = \mathcal{T} \mathcal{J}_R(\xi) = \mathcal{J}_R(-\xi)$$

# Pose Lie Groups and Lie Algebras

	Lie algebra		Lie group
$4 \times 4$	$\xi^\wedge \in \mathfrak{se}(3)$	$\xrightarrow{\exp}$	$\mathbf{T} \in SE(3)$
	$\downarrow \text{ad}$		$\downarrow \text{Ad}$
$6 \times 6$	$\xi^\wedge \in \text{ad}(\mathfrak{se}(3))$	$\xrightarrow{\exp}$	$\mathcal{T} \in \text{Ad}(SE(3))$

---


$$\begin{aligned}
 \mathcal{T} &= \underbrace{\text{Ad} \left( \exp(\hat{\xi}) \right)}_{\mathcal{T}} = \exp \left( \underbrace{\text{ad}(\hat{\xi})}_{\xi^\wedge} \right) \quad \xi = \begin{bmatrix} \rho \\ \theta \end{bmatrix} \in \mathbb{R}^6 \\
 &= \text{Ad} \left( \exp \left( \begin{bmatrix} \hat{\theta} & \rho \\ \mathbf{0}^T & 0 \end{bmatrix} \right) \right) = \exp \left( \text{ad} \left( \begin{bmatrix} \hat{\theta} & \rho \\ \mathbf{0}^T & 0 \end{bmatrix} \right) \right) \\
 &= \text{Ad} \left( \begin{bmatrix} \exp(\hat{\theta}) & J_L(\theta)\rho \\ \mathbf{0}^T & 1 \end{bmatrix} \right) = \exp \left( \begin{bmatrix} \hat{\theta} & \hat{\rho} \\ \mathbf{0} & \hat{\theta} \end{bmatrix} \right) \\
 &= \begin{bmatrix} \exp(\hat{\theta}) & (J_L(\theta)\rho)^\wedge \exp(\hat{\theta}) \\ \mathbf{0} & \exp(\hat{\theta}) \end{bmatrix}
 \end{aligned}$$

## Rodrigues Formula for the Adjoint of $SE(3)$

- ▶ The exponential map is **surjective** but **not injective**, i.e., every element of  $Ad(SE(3))$  can be generated from multiple elements of  $ad(\mathfrak{se}(3))$
- ▶ **Rodrigues Formula**: using  $(\overset{\wedge}{\xi})^5 + 2\|\theta\|^2(\overset{\wedge}{\xi})^3 + \|\theta\|^4\overset{\wedge}{\xi} = 0$  we can obtain a direct expression of  $\mathcal{T} \in Ad(SE(3))$  in terms of  $\xi = \begin{bmatrix} \rho \\ \theta \end{bmatrix} \in \mathbb{R}^6$ :

$$\begin{aligned}\mathcal{T} = Ad(T) &= \exp\left(\overset{\wedge}{\xi}\right) = \begin{bmatrix} \exp(\hat{\theta}) & (J_L(\theta)\rho)^\wedge \exp(\hat{\theta}) \\ \mathbf{0} & \exp(\hat{\theta}) \end{bmatrix} = \sum_{n=0}^{\infty} \frac{1}{n!} (\overset{\wedge}{\xi})^n \\ &= I + \left( \frac{3 \sin \|\theta\| - \|\theta\| \cos \|\theta\|}{2\|\theta\|} \right) \overset{\wedge}{\xi} + \left( \frac{4 - \|\theta\| \sin \|\theta\| - 4 \cos \|\theta\|}{2\|\theta\|^2} \right) (\overset{\wedge}{\xi})^2 \\ &\quad + \left( \frac{\sin \|\theta\| - \|\theta\| \cos \|\theta\|}{2\|\theta\|^3} \right) (\overset{\wedge}{\xi})^3 + \left( \frac{2 - \|\theta\| \sin \|\theta\| - 2 \cos \|\theta\|}{2\|\theta\|^4} \right) (\overset{\wedge}{\xi})^4\end{aligned}$$

## Distances in $SE(3)$

- Two ways to define differences between  $SE(3)$  and  $Ad(SE(3))$  elements:

$$\xi_{12} = \log (T_1^{-1} T_2)^{\vee} = \log (\mathcal{T}_1^{-1} \mathcal{T}_2)^{\gamma}$$

$$\xi_{21} = \log (T_2 T_1^{-1})^{\vee} = \log (\mathcal{T}_2 \mathcal{T}_1^{-1})^{\gamma}$$

- Inner product on  $\mathfrak{se}(3)$  and  $ad(\mathfrak{se}(3))$ :

$$\langle \hat{\xi}_1, \hat{\xi}_2 \rangle = \text{tr} \left( \hat{\xi}_1 \begin{bmatrix} \frac{1}{2}I & \mathbf{0} \\ \mathbf{0}^\top & 1 \end{bmatrix} \hat{\xi}_2^\top \right) = \xi_1^\top \xi_2$$

$$\langle \overset{\wedge}{\xi}_1, \overset{\wedge}{\xi}_2 \rangle = \text{tr} \left( \overset{\wedge}{\xi}_1 \begin{bmatrix} \frac{1}{4}I & \mathbf{0} \\ \mathbf{0} & \frac{1}{2}I \end{bmatrix} \overset{\wedge}{\xi}_2^\top \right) = \xi_1^\top \xi_2$$

- The right and left distances on  $SE(3)$  and  $Ad(SE(3))$  are:

$$\xi_{12} = \sqrt{\langle \hat{\xi}_{12}, \hat{\xi}_{12} \rangle} = \sqrt{\langle \overset{\wedge}{\xi}_{12}, \overset{\wedge}{\xi}_{12} \rangle} = \sqrt{\xi_{12}^\top \xi_{12}} = \|\xi_{12}\|_2$$

$$\xi_{21} = \sqrt{\langle \hat{\xi}_{21}, \hat{\xi}_{21} \rangle} = \sqrt{\langle \overset{\wedge}{\xi}_{21}, \overset{\wedge}{\xi}_{21} \rangle} = \sqrt{\xi_{21}^\top \xi_{21}} = \|\xi_{21}\|_2$$

## Integration in $SE(3)$

- ▶ The distance between a pose  $T = \exp(\hat{\xi})$  and a small perturbation  $\exp((\xi + \delta\xi)^\wedge)$  can be approximated using the BCH formulas:

$$\log \left( \exp(\hat{\xi})^{-1} \exp((\xi + \delta\xi)^\wedge) \right)^\vee \approx \mathcal{J}_R(\xi) \delta\xi$$

$$\log \left( \exp((\xi + \delta\xi)^\wedge) \exp(\hat{\xi})^{-1} \right)^\vee \approx \mathcal{J}_L(\xi) \delta\xi$$

- ▶ Regardless whether the left or the right distance metric is used, the infinitesimal volume element is:

$$|\det(\mathcal{J}(\xi))| = |\det(J(\theta))|^2 = 4 \left( \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right)^2$$

- ▶ Integrating functions of poses can then be carried out as follows:

$$\int_{SE(3)} f(T) dT = \int_{\mathbb{R}^3, \|\theta\| < \pi} f \left( \exp(\hat{\xi}) \right) |\det(\mathcal{J}(\xi))| d\xi$$



## Lie Algebra $\mathfrak{se}(3)$ Identities

$$\hat{\xi} = \begin{bmatrix} \hat{\rho} \\ \hat{\theta} \end{bmatrix} = \begin{bmatrix} \hat{\theta} & \rho \\ \mathbf{0}^\top & 0 \end{bmatrix} \in \mathbb{R}^{4 \times 4} \quad \overset{\wedge}{\xi} = ad(\hat{\xi}) = \begin{bmatrix} \overset{\wedge}{\rho} \\ \hat{\theta} \end{bmatrix} = \begin{bmatrix} \hat{\theta} & \hat{\rho} \\ \mathbf{0} & \hat{\theta} \end{bmatrix} \in \mathbb{R}^{6 \times 6}$$

$$\overset{\wedge}{\zeta} \overset{\wedge}{\xi} = -\overset{\wedge}{\xi} \zeta \quad \zeta \in \mathbb{R}^6$$

$$\overset{\wedge}{\xi} \overset{\wedge}{\xi} = 0$$

$$\hat{\xi}^4 + (\mathbf{s}^\top \mathbf{s}) \hat{\xi}^2 = 0 \quad \mathbf{s} \in \mathbb{R}^3$$

$$\left(\overset{\wedge}{\xi}\right)^5 + 2(\mathbf{s}^\top \mathbf{s}) \left(\overset{\wedge}{\xi}\right)^3 + (\mathbf{s}^\top \mathbf{s})^2 \overset{\wedge}{\xi} = 0$$

$$\mathbf{m}^\odot := \begin{bmatrix} \mathbf{s} \\ \lambda \end{bmatrix}^\odot = \begin{bmatrix} \lambda I & -\hat{\mathbf{s}} \\ \mathbf{0}^\top & \mathbf{0}^\top \end{bmatrix} \in \mathbb{R}^{4 \times 6} \quad \mathbf{m}^\odot := \begin{bmatrix} \mathbf{s} \\ \lambda \end{bmatrix}^\odot = \begin{bmatrix} \mathbf{0} & \mathbf{s} \\ -\hat{\mathbf{s}} & \mathbf{0} \end{bmatrix} \in \mathbb{R}^{6 \times 4}$$

$$\hat{\xi} \mathbf{m} = \mathbf{m}^\odot \xi \quad \mathbf{m}^\top \hat{\xi} = \xi^\top \mathbf{m}^\odot$$

## Lie Group $SE(3)$ Identities

$$T = \exp(\hat{\xi}) = \begin{bmatrix} \exp(\hat{\theta}) & J_L(\theta)\rho \\ \mathbf{0}^T & 1 \end{bmatrix}$$
$$= \sum_{n=0}^{\infty} \frac{1}{n!} \hat{\xi}^n = I + \hat{\xi} + \left( \frac{1 - \cos \|\theta\|}{\|\theta\|^2} \right) \hat{\xi}^2 + \left( \frac{\|\theta\| - \sin \|\theta\|}{\|\theta\|^3} \right) \hat{\xi}^3 \approx I + \hat{\xi}$$

$$T^{-1} = \exp(-\hat{\xi}) = \begin{bmatrix} \exp(-\hat{\theta}) & -\exp(-\hat{\theta}) J_L(\theta)\rho \\ \mathbf{0}^T & 1 \end{bmatrix} = \sum_{n=0}^{\infty} \frac{1}{n!} (-\hat{\xi})^n \approx I - \hat{\xi}$$

$$\det(T) = 1$$

$$\text{tr}(T) = 2 \cos \|\theta\| + 2$$

$$T \hat{\xi} = \hat{\xi} T$$

# Lie Group $Ad(SE(3))$ Identities

$$\begin{aligned}\mathcal{T} = Ad(T) &= \exp \left( \overset{\wedge}{\xi} \right) = \begin{bmatrix} \exp(\hat{\theta}) & (J_L(\theta)\rho)^\wedge \exp(\hat{\theta}) \\ \mathbf{0} & \exp(\hat{\theta}) \end{bmatrix} \\ &= \sum_{n=0}^{\infty} \frac{1}{n!} \overset{\wedge}{\xi}^n = I + \left( \frac{3 \sin \|\theta\| - \|\theta\| \cos \|\theta\|}{2\|\theta\|} \right) \overset{\wedge}{\xi} + \left( \frac{4 - \|\theta\| \sin \|\theta\| - 4 \cos \|\theta\|}{2\|\theta\|^2} \right) (\overset{\wedge}{\xi})^2 \\ &\quad + \left( \frac{\sin \|\theta\| - \|\theta\| \cos \|\theta\|}{2\|\theta\|^3} \right) (\overset{\wedge}{\xi})^3 + \left( \frac{2 - \|\theta\| \sin \|\theta\| - 2 \cos \|\theta\|}{2\|\theta\|^4} \right) (\overset{\wedge}{\xi})^4 \approx I + \overset{\wedge}{\xi}\end{aligned}$$

$$\mathcal{T}^{-1} = \exp \left( -\overset{\wedge}{\xi} \right) = \begin{bmatrix} \exp(-\hat{\theta}) & -\exp(-\hat{\theta}) (J_L(\theta)\rho)^\wedge \\ \mathbf{0} & \exp(-\hat{\theta}) \end{bmatrix} = \sum_{n=0}^{\infty} \frac{1}{n!} \left( -\overset{\wedge}{\xi} \right)^n \approx I - \overset{\wedge}{\xi}$$

$$\mathcal{T}\xi = \xi$$

$$\overset{\wedge}{\mathcal{T}}\overset{\wedge}{\xi} = \overset{\wedge}{\xi}\overset{\wedge}{\mathcal{T}}$$

$$(\mathcal{T}\zeta)^\wedge = \mathcal{T}\hat{\zeta}\mathcal{T}^{-1}$$

$$(\overset{\wedge}{\mathcal{T}}\overset{\wedge}{\zeta}) = \overset{\wedge}{\mathcal{T}}\overset{\wedge}{\zeta}\overset{\wedge}{\mathcal{T}}^{-1} \quad \zeta \in \mathbb{R}^6$$

$$\exp((\mathcal{T}\zeta)^\wedge) = \mathcal{T} \exp(\hat{\zeta}) \mathcal{T}^{-1}$$

$$\exp\left(\overset{\wedge}{(\mathcal{T}\zeta)}\right) = \overset{\wedge}{\mathcal{T}} \exp\left(\overset{\wedge}{\zeta}\right) \overset{\wedge}{\mathcal{T}}^{-1}$$

$$(T\mathbf{m})^\odot = T\mathbf{m}^\odot\mathcal{T}^{-1}$$

$$((T\mathbf{m})^\odot)^T (T\mathbf{m})^\odot = \mathcal{T}^{-T} (\mathbf{m}^\odot)^T \mathbf{m}^\odot \mathcal{T}^{-1}$$

## $SO(3)$ and $SE(3)$ Kinematics

## Rotation Kinematics

- ▶ The trajectory  $R(t)$  of a continuous rotation motion should satisfy:

$$R^\top(t)R(t) = I \quad \Rightarrow \quad \dot{R}^\top(t)R(t) + R^\top(t)\dot{R}(t) = 0.$$

- ▶ The matrix  $R^\top(t)\dot{R}(t)$  is **skew-symmetric** and there must exist some vector-valued function  $\omega(t) \in \mathbb{R}^3$  such that:

$$R^\top(t)\dot{R}(t) = \hat{\omega}(t) \quad \Rightarrow \quad \boxed{\dot{R}(t) = R(t)\hat{\omega}(t)}$$

- ▶ A skew-symmetric matrix gives a first order approximation to a rotation matrix:

$$R(t + dt) \approx R(t) + R(t)\hat{\omega}(t)dt$$

## Rotation Kinematics

- ▶ Let  $R \in SO(3)$  be the orientation of a rigid body rotating with angular velocity  $\omega \in \mathbb{R}^3$  with respect to the world frame.
- ▶ **Rotation kinematic equations of motion:**

$$\dot{R} = R\hat{\omega}_B = \hat{\omega}_W R$$

where  $\omega_B$  and  $\omega_W := R\omega_B$  are the body-frame and world-frame coordinates of  $\omega$ , respectively.

- ▶ Assuming  $\omega$  is constant over a short period  $\tau$ :

$$R(t + \tau) = R(t) \exp(\tau \hat{\omega}_B) = \exp(\tau \hat{\omega}_W) R(t)$$

- ▶ **Discrete Rotation Kinematics:** let  $R_k := R(t_k)$ ,  $\tau_k := t_{k+1} - t_k$ , and  $\omega_k := \omega_B(t_k)$  leading to:

$$R_{k+1} = R_k \exp(\tau_k \hat{\omega}_k)$$

## Pose Kinematics

► **Angular velocity:**  $R^\top(t)\dot{R}(t) = I \Rightarrow R^\top(t)\dot{R}(t) = \hat{\omega}(t) \in \mathfrak{so}(3)$

► **Twist:** similarly for  $T(t) \in SE(3)$  consider:

$$T^{-1}(t)\dot{T}(t) = \begin{bmatrix} R^\top(t)\dot{R}(t) & R^\top(t)\dot{\mathbf{p}}(t) \\ 0 & 0 \end{bmatrix} = \begin{bmatrix} \hat{\omega}(t) & \mathbf{v}(t) \\ 0 & 0 \end{bmatrix} \in \mathfrak{se}(3)$$

where  $\hat{\omega}(t) := R^\top(t)\dot{R}(t)$  and  $\mathbf{v}(t) := R^\top(t)\dot{\mathbf{p}}(t)$  are the **body-frame** angular and linear velocities of the body

► **Generalized velocity:**  $\zeta(t) := \begin{bmatrix} \mathbf{v}(t) \\ \omega(t) \end{bmatrix} \in \mathbb{R}^6$

►  $\zeta(t)$  is the velocity of the body frame moving relative to the world frame as viewed in the **body frame**

► **Continuous-time Pose Kinematics:**  $\dot{T}(t) = T(t)\hat{\zeta}(t)$

► **Discrete-time Pose Kinematics:**  $T_{k+1} = T_k \exp\left(\tau_k \hat{\zeta}_k\right)$

# Pose Kinematics

- ▶ Consider a moving body frame  $\{B\}$  with pose  $T(t) \in SE(3)$
- ▶ Let  $\mathbf{s}_B \in \mathbb{R}^3$  be a point in the body frame with homogeneous coordinates  $\underline{\mathbf{s}}_B$
- ▶ The velocity of  $\mathbf{s}_B$  with respect to the world frame  $\{W\}$  can be determined as follows:

$$\begin{aligned}\underline{\mathbf{s}}_W(t) &= T(t)\underline{\mathbf{s}}_B \\ \dot{\underline{\mathbf{s}}}_W(t) &= \dot{T}(t)\underline{\mathbf{s}}_B = \dot{T}(t)T(t)^{-1}\underline{\mathbf{s}}_W(t) \\ &= T(t)\hat{\zeta}(t)T(t)^{-1}\underline{\mathbf{s}}_W(t) \\ &= \begin{bmatrix} R(t)\hat{\omega}(t)R(t)^\top & R(t)\mathbf{v}(t) - R(t)\hat{\omega}(t)R(t)^\top \mathbf{p}(t) \\ \mathbf{0}^\top & 0 \end{bmatrix} \begin{bmatrix} \mathbf{s}_W(t) \\ 1 \end{bmatrix} \\ &= \begin{bmatrix} (R(t)\omega(t))^\wedge (\mathbf{s}_W(t) - \mathbf{p}(t)) + R(t)\mathbf{v}(t) \\ 1 \end{bmatrix}\end{aligned}$$